Isoneutral Mixing Operators In Oceanic Models : Numerical Delicacies

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Grenoble



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Numerical Models

Local mesh refinement, ocean-atmosphere coupling, numerical methods

E. Blayo, L. Debreu, F. Lemarié

Variational Data assimilation

F.-X. LeDimet, A. Vidard, E. Kazantsev, M. Nodet

Uncertainty quantification

C. Prieur, C. Helbert

- Spurious diapycnal mixing: The problem and a solution, Marchesiello et al (Ocean Modelling, 2009)
- 2 On the stability and accuracy of the harmonic and biharmonic isoneutral mixing operators in ocean models, Lemarié et al (Ocean Modelling, 2012)
- 3 Other problems and perspectives, Demange et al (in preparation)

- Spurious diapycnal mixing: The problem and a solution, Marchesiello et al (Ocean Modelling, 2009)
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Spurious diapycnal mixing: The problem and a solution

• Tracer mixing in a stratified ocean (away from turbulent boundary layers)

- $\kappa_{\rm dia} \approx 10^{-5} {\rm m s}^{-1}$ (e.g. Ledwell et al., 1993)
- $\kappa_{
 m iso} pprox 10^3 {
 m m s}^{-1}$ (horizontal scale pprox 100 km)

Vertical coordinates systems:

Main problems of terrain following (σ) (and geopotential z) models: Pressure gradient error and Diapycnal Mixing Stronger in ocean models than in atmospheric models



Spurious diapycnal mixing: The problem and a solution

ROMS model, terrain following coordinates, Third order upstream biased scheme for tracers. Salinity at 1000m depth



The (implicit) diffusion of the upstream biased scheme acts along horizontal coordinates . . .

ROMS model, σ coordinates, Third order upwind scheme for tracers Salitiny at 1000m after 2 years of integration



A solution

Rotate the diffuse part of the upstream biased scheme:



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Rotation of the diffusive part introduces mixed (horizontal/vertical) and purely vertical derivatives

Stabilization of a (high order) rotated diffusion operator

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Isoneutral mixing problem (continuous formulation)

Isoneutral Laplacian operator

Tracer field q in an unbounded domain $\Omega = \mathbb{R}^3$

$$\begin{cases} \partial_t q &= \mathcal{D}_2(q) = \nabla \cdot (\mathbf{R} \nabla q) = -\nabla \cdot \mathbf{F} & \text{in } \Omega \times [0, T], \\ q|_{t=0} &= q_0(x, y, z) & \text{in } \Omega, \end{cases}$$

the matrix form of the tensor ${\boldsymbol{\mathsf{R}}}$ is

[Redi, 1982; Gent & McWilliams, 1990]

$$\mathbf{R} = \begin{pmatrix} \kappa_x & 0 & \kappa_x \alpha_x \\ 0 & \kappa_y & \kappa_y \alpha_y \\ \kappa_x \alpha_x & \kappa_y \alpha_y & \kappa_x \alpha_x^2 + \kappa_y \alpha_y^2 \end{pmatrix}, \quad \text{with} \quad \boldsymbol{\alpha} = (\alpha_x, \alpha_y) = -\frac{(\partial_x \rho, \partial_y \rho)}{\partial_z \rho}, \quad \|\boldsymbol{\alpha}\| \ll 1.$$

Properties

- Orthogonality condition : $\mathbf{F} \cdot \boldsymbol{\rho}_{\perp} = 0$
- Satisfy monotonicity principle [Mathieu and Deleersnijder, 1998]
- Satisfy global tracer variance dissipation [Griffies et al., 1998]

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Isoneutral mixing problem (continuous formulation)

Isoneutral biharmonic operator

The rotated biharmonic operator reads

$$\begin{cases} \partial_t q &= \mathcal{D}_4(q) = -\mathcal{D}_2(\Psi) = -\nabla \cdot \mathbf{F_4}, \quad \text{with } \Psi = \mathcal{D}_2(q), \\ q|_{t=0} &= q_0(x, y, z), \end{cases}$$

composition of two isoneutral Laplacian operators

Properties

- Orthogonality condition : $\mathbf{F_4} \cdot \boldsymbol{
 ho}_\perp = 0$
- Satisfy global tracer variance dissipation [Griffies, 2004]

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Time discretization

[Lemarié et al., 2012]

Objectives :

- same stability limit of the non-rotated operators
- only the vertical direction can be implicit

Proposed scheme :

 \rightarrow Laplacian

$$\mathcal{D}_{2}(q) = \underbrace{\partial_{x}(\kappa_{x} \left[\partial_{x}q + \alpha_{x}\partial_{z}q\right]) + \partial_{z}\left(\alpha_{x}\kappa_{x}\partial_{x}q\right)}_{explicit} + \underbrace{\partial_{z}\left(\kappa_{x}\alpha_{x}^{2}\partial_{z}q\right)}_{G_{3}(q)}$$

$$\begin{cases} q^{\star} = q^{n} + \Delta t \mathcal{D}_{2}(q^{n}) \\ q^{n+1} = q^{\star} + \theta \Delta t \left[G_{3}(q^{n+1}) - G_{3}(q^{n})\right] \end{cases}$$

Method of Stabilizing Corrections [van der Houwen & Verwer, 1979; Hundsdorfer, 2002]

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Time discretization

[Lemarié et al., 2012]

Objectives :

- same stability limit of the non-rotated operators
- only the vertical direction can be implicit

Proposed scheme :

 \rightarrow Biharmonic

$$\left\{ egin{array}{rcl} q^{\star} &=& q^n + \Delta t \mathcal{D}_4(q^n) \ \\ q^{n+1} &=& q^{\star} + \Delta t \partial_z \left[\widetilde{\kappa} \partial_z q^{n+1} - \widetilde{\kappa} \partial_z q^n
ight] \end{array}
ight.$$

 $\widetilde{\kappa}$ chosen through linear stability analysis

$$\widetilde{\kappa} = 8\Delta z^2 \sigma_x s_x^2 (1+s_x^2)/\Delta t \qquad \Rightarrow \qquad \sigma_x \leq rac{1}{8}$$

with $s_x = \alpha_x \frac{\Delta x}{\Delta z}$, $\sigma_x = \kappa_x \frac{\Delta t}{\Delta x^4}$ Debreu et al. (Inria) Isoneutral Mixing Operators In Oceanic Mod. April, 11, 2014 14 / 23

Conclusions of Lemarié et al (2012)

- \triangleright the rotated operators can be advanced with the same time step as the non-rotated ones !!
- > the rotated biharmonic can be a viable operator for use in high-resolution global models
- \triangleright clipping/tapering should act on **s** rather than lpha
- arepsilon a slope-dependent discretization provides more accuracy notably for large grid slope ratios
- ▷ has been implemented in ROMS and NEMO (also for rotation of hyperviscosity)

Alternatives to the use of rotated operators:

- No major improvements so far with higher-order advection schemes ...
- Arbitrary-Lagrangian-Eulerian (ALE) vertical coordinate \rightarrow strong ongoing effort

Remaining open issue:

Monotone advection schemes and diapycnal mixing

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- Amount of numerical viscosity and vertical advection schemes
- Barotropic/Baroclinic time splitting
- Internal Gravity Wave propagation

Other sources of artificial diffusion (and dispersion)

• Amount of numerical viscosity and vertical advection schemes



Initial condition

Momentum : UP3

Momentum : UP1

- Barotropic/Baroclinic time splitting
- Internal Gravity Wave propagation

Other sources of artificial diffusion (and dispersion)

- Amount of numerical viscosity and vertical advection schemes
- Barotropic/Baroclinic time splitting



- Amount of numerical viscosity and vertical advection schemes
- Barotropic/Baroclinic time splitting
- Internal Gravity Wave propagation

$$\frac{\partial u}{\partial t} + \frac{1}{\rho_0} \frac{\partial p}{\partial x} = 0$$
$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0$$
$$\frac{\partial p}{\partial z} = -\rho g$$
$$\frac{\partial p}{\partial t} + w \frac{\partial \bar{p}}{\partial z} = 0$$

Very basic numerics:

- Second order discretization
- Computational mode of the Lorenz grid
- No dissipation

Decomposition into vertical modes and characteristic variables

$$u(x, z, t) = \sum u_n(x, t)M_n(z)$$

$$\rho(x, z, t) = -\rho_0 \sum h_n(x, t) \frac{\mathrm{d}M_n(z)}{\mathrm{d}z}$$

$$\frac{\partial}{\partial t}(u_n \pm \frac{g}{c_n}h_n) \pm c_n \frac{\partial}{\partial x}(u_n \pm \frac{g}{c_n}h_n) = 0$$

Other sources of artificial diffusion (and dispersion)

- Amount of numerical viscosity and vertical advection schemes
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Monotonicity of characteristic variables in the primitive variables formulation ?

$$\begin{cases} \frac{\partial u}{\partial t} + \frac{1}{\rho_0} \frac{\partial p}{\partial x} = \frac{c_1 \Delta x}{2} \frac{\partial^2 u}{\partial x^2} \\ \frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0 \\ \frac{\partial p}{\partial z} = -\rho g \\ \frac{\partial \rho}{\partial t} + w \frac{\partial \bar{p}}{\partial z} = \frac{c_1 \Delta x}{2} \frac{\partial^2 \rho}{\partial x^2} \\ \frac{\partial \rho}{\partial x^2}$$

Internal Wave propagation



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- Marchesiello, P., Debreu, L., Couvelard, X., 2009. Spurious diapycnal mixing in terrain-following coordinate models: The problem and a solution. Ocean Modell. 26, 159-169.
- Lemarié, F., Debreu, L., Shchepetkin, A.F., McWilliams, J.C., 2012. On the stability and accuracy of the harmonic and biharmonic isoneutral mixing operators in ocean models, Ocean Modell., 52-53, 9-35.
- Demange J., Debreu L., Marchesiello P., 2014: On the use of a depth-dependent barotropic mode for free surface ocean models, to be submitted
- Demange J., Debreu L., Marchesiello P., 2014: Vertical mode decomposition and dissipation, to be submitted

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Properties of the discretized operator ?

$$\mathcal{D}_{2}(q) = \partial_{x} \left(\kappa_{x} \frac{\mathcal{J}_{x}(q,\rho)}{\partial_{z}\rho} \right) - \partial_{z} \left(\kappa_{x} \frac{\partial_{x}\rho}{\partial_{z}\rho} \frac{\mathcal{J}_{x}(q,\rho)}{\partial_{z}\rho} \right)$$

 $\mathcal{J}_{\mathsf{x}}^{\mathsf{u}} = (\partial_{\mathsf{x}} q) \left(\overline{\partial_{\mathsf{z}} \rho} \right) - (\partial_{\mathsf{x}} \rho) \left(\overline{\partial_{\mathsf{z}} q} \right), \qquad \mathcal{J}_{\mathsf{x}}^{\mathsf{w}} = \left(\overline{\partial_{\mathsf{x}} q} \right) \left(\partial_{\mathsf{z}} \rho \right) - \left(\overline{\partial_{\mathsf{x}} \rho} \right) \left(\partial_{\mathsf{z}} q \right).$



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Properties of the discretized operator ?

$$\mathcal{D}_{2}(q) = \partial_{x} \left(\kappa_{x} \frac{\mathcal{J}_{x}(q,\rho)}{\partial_{z}\rho} \right) - \partial_{z} \left(\kappa_{x} \frac{\partial_{x}\rho}{\partial_{z}\rho} \frac{\mathcal{J}_{x}(q,\rho)}{\partial_{z}\rho} \right)$$

$$\mathcal{J}_{x}^{u} = (\partial_{x}q) \, \widetilde{(\partial_{z}\rho)} - (\partial_{x}\rho) \, \widetilde{(\partial_{z}q)}, \qquad \mathcal{J}_{x}^{w} = \widetilde{(\partial_{x}q)} \, (\partial_{z}\rho) - \widetilde{(\partial_{x}\rho)} \, (\partial_{z}q) \, .$$

	ref.	min-max	TVD ¹	$\mathbf{F}(ho)=0$	misc.
COX	[Cox, 1987]	no	no	no	$2\Delta x \mod x$
TRIADS	[Griffies et al., 1998]	no	yes	yes	
SW-TRIADS	[Lemarié et al., 2012]	no	yes	yes	most compact
					stencil

 \rightarrow all linear schemes produce over/under shootings [Beckers et al., 2000]

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¹Tracer Variance Dissipation

Spatial discretization and stability limits

TRIADS

	TRIADS	SW-TRIADS	
Laplacian			
Rotated	$\sigma_x(1+s_x^2) \leq 1/2$	$\sigma_x \max(1, s_x^2) \leq 1/2$	
Non-rotated	$\sigma_x \leq 1/2$		
Biharmonic			
Rotated	$\left(\sigma_x^{(4)}(1+s_x^2) ight)^2\leq 1/8$	$\left(\sigma_x^{(4)}\max(1,s_x^2) ight)^2\leq 1/8$	
Non-rotated	$\left(\sigma_{x}^{(4)} ight)^{2}\leq1/8$		
with $\sigma_x = \kappa_x \frac{\Delta}{\Delta x}$	$\frac{t}{z^2}$, $s_x = \alpha_x \frac{\Delta x}{\Delta z}$,	$\sigma_x^{(4)} = \sqrt{B_x} \frac{\sqrt{\Delta t}}{\Delta x^2}$	
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Adaptive Mesh Refinement for ocean models



Blayo and Debreu, 1999: Adaptive Mesh Refinement for Finite-Difference Ocean Models: First experiments, JPO

Ptp. 3. Numerical solution of the modon problem at three different time steps (CI: 1000 m2 s-1) for the (a) nested-grid calculation and (b) AMR calculation. The location of the fine erid is indicated



Debreu et al, 2005: A general adaptive multi-resolution approach to ocean modelling : experiments in a primitive equation model of the north atlantic. Adaptive Mesh Refinement - Theory and Applications, Lecture Notes in Computational Science and Engineering



Fig. 3. Barotropic streamfunction : from left to right: uniform 1/3° simulation (reference simulation), uniform 1/2° simulation, uniform uniform 1° simulation, adaptive 1°-1/3° simulation. Contour interval is of 5 Sverdrup

- \rightarrow refinement criterion ?
- \rightarrow How to define the bathymetry across the grid hierarchy ?

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Adaptive Mesh Refinement for ocean models

Before doing adaptive mesh refinement, are we sure than fixed mesh refinement is doing well ?

- Debreu and Blayo, 2008: Two-way embedding algorithms : a review. Ocean Dynamics
- Debreu et al, 2012: Two-way nesting in split-explicit ocean models : Algorithms, implementation and validation. Ocean Modelling
 - High order update schemes / conservation and refluxing
 - Coupling at the fast modes (barotropic modes) level
- Debreu et al, 2008: AGRIF: Adaptive grid refinement in fortran. Comput. Geosci. A library implementing Berger and Oliger algorithm and a source-to-source (monogrid to multigrid) translator

Current developments in ocean models:

Different vertical grids and potentially different vertical coordinate systems between coarse and fine grids