



A Discretization of Deep Atmospheric Model Dynamics for NCEP Global Forecast System

Hann-Ming Henry Juang

Environment Modeling Center, NOAA/NWS/NCEP, Washington, DC

Introduction

- While NCEP GFS extends its vertical domain to couple with space environmental model, we called this version of GFS as Whole Atmospheric Model (WAM). WAM is a hydrostatic system with enthalpy as thermodynamic variable (Juang 2011 MWR).
- We propose to do **deep atmospheric** dynamics for NCEP GFS to support WAM
- Here, we would like to re-iterate the reasons to use deep atmospheric dynamics and to illustrate the discretization of deep atmospheric dynamics for NCEP GFS.

Opr GFS

vs

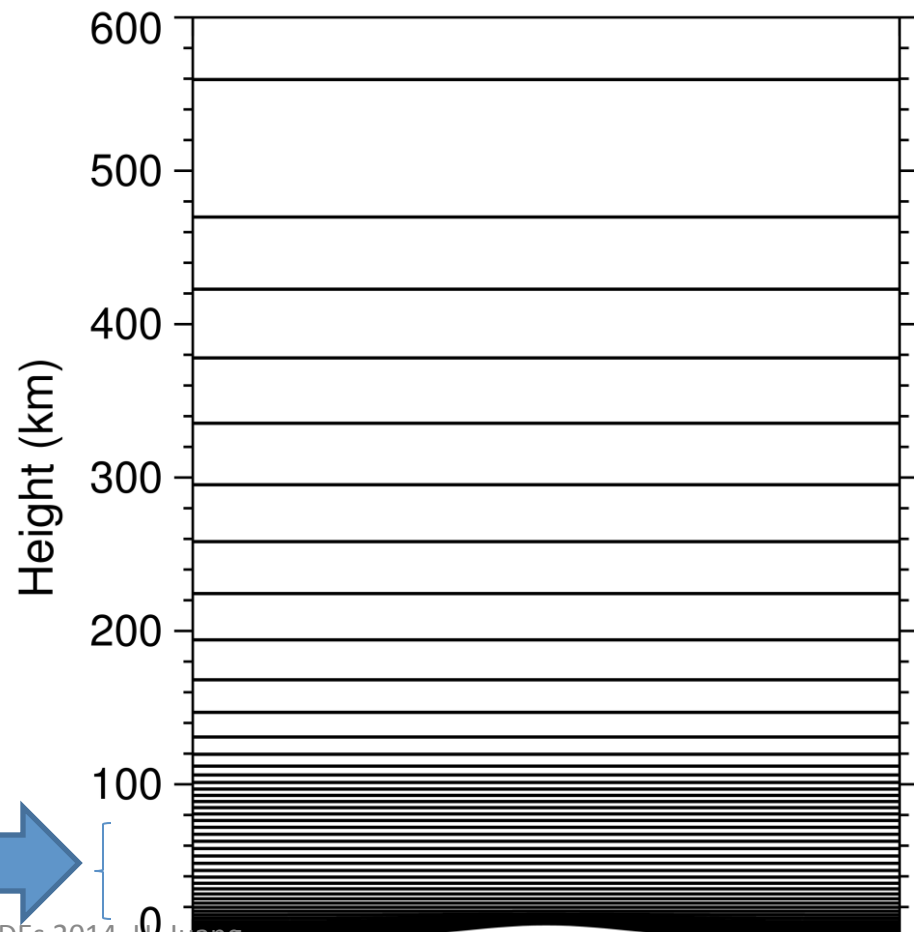
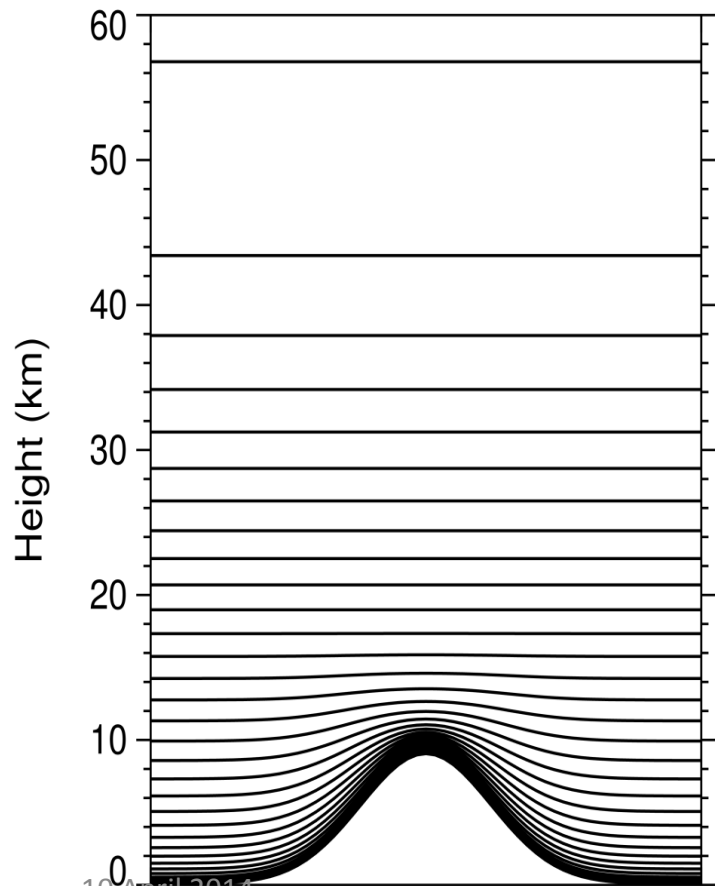
WAM

64 layers

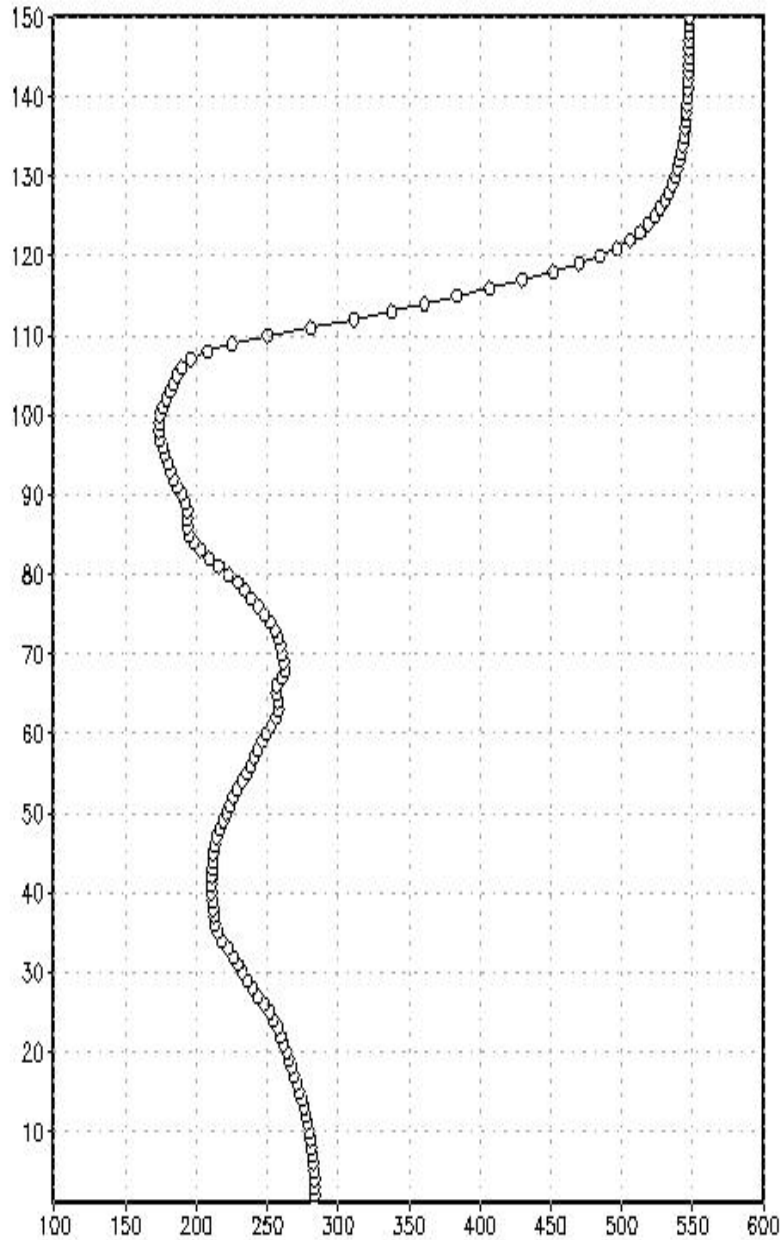
GFS hybrid vertical grid
(every 2nd level)

150 layers

WAM hybrid vertical grid
(every 3rd level)



WAM T at lat=29.5 lon=0



Example of T profile of 150 layers

WAM uses generalized hybrid coordinate with enthalpy $C_p T$ as thermodynamics variables, where C_p is summation of each gases.

	R	C_p
O	519.674	1299.18
O2	259.837	918.096
O3	173.225	820.239
Dry air	296.803	1039.64
H2O	461.50	1846.00

Maxima wind (m/s) at NCEP GFS 150 layers WAM

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in do_dynamics_two_loop for spdmx at kdt= 40825
spdmx(001:010)= 19. 20. 21. 23. 25. 26. 27. 28. 28. 28.
spdmx(011:020)= 28. 27. 27. 27. 27. 28. 28. 28. 29. 30.
spdmx(021:030)= 31. 33. 35. 37. 40. 42. 44. 46. 49. 53.
spdmx(031:040)= 58. 61. 63. 63. 62. 60. 55. 47. 45. 44.
spdmx(041:050)= 45. 45. 47. 49. 52. 55. 59. 62. 65. 68.
spdmx(051:060)= 72. 76. 80. 84. 87. 90. 93. 95. 97. 98.
spdmx(061:070)= 102. 110. 118. 127. 135. 143. 149. 153. 155. 152.
spdmx(071:080)= 147. 145. 142. 138. 135. 132. 130. 126. 121. 119.
spdmx(081:090)= 114. 112. 110. 106. 100. 95. 94. 90. 89. 89.
spdmx(091:100)= 87. 82. 91. 95. 99. 97. 104. 100. 111. 120.
spdmx(101:110)= 125. 133. 148. 167. 172. 164. 159. 160. 147. 124.
spdmx(111:120)= 117. 125. 133. 138. 137. 157. 183. 202. 220. 243.
spdmx(121:130)= 269. 297. 319. 338. 355. 368. 378. 386. 392. 396.
spdmx(131:140)= 399. 402. 404. 405. 406. 407. 408. 409. 410. 410.
spdmx(141:150)= 411. 412. 412. 413. 413. 414. 414. 415. 415. 418.
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Shallow ($r=a$) vs Deep ($r=a+z$)

- Assume at $z=637.12\text{km}$, so $r = 1.1a$

- For shallow dynamic

$$u = a \cos \phi \frac{d\lambda}{dt}$$

- For deep atmosphere

$$u = r \cos \phi \frac{d\lambda}{dt}$$

- For example, $\phi=45^\circ$, $u=400\text{m/s}$, after one hour advection, the displacement has about 1° error in λ .

Deep atmospheric equation in height & spherical coordinates

Momentum

$$\frac{du}{dt} - \frac{uv \tan \phi}{r} + \frac{uw}{r} - (2\Omega \sin \phi)v + (2\Omega \cos \phi)w + \frac{1}{\rho r \cos \phi} \frac{\partial p}{\partial \lambda} = F_u$$

$$\frac{dv}{dt} + \frac{u^2 \tan \phi}{r} + \frac{vw}{r} + (2\Omega \sin \phi)u + \frac{1}{\rho r} \frac{\partial p}{\partial \phi} = F_v$$

$$\frac{dw}{dt} - \frac{u^2 + v^2}{r} - (2\Omega \cos \phi)u + \frac{1}{\rho} \frac{\partial p}{\partial r} + g = F_w$$

where

$$\frac{dA}{dt} = \frac{\partial A}{\partial t} + u \frac{\partial A}{r \cos \phi \partial \lambda} + v \frac{\partial A}{r \partial \phi} + w \frac{\partial A}{\partial r} \quad r = a + z$$

$$u = r \cos \phi \frac{d\lambda}{dt}$$

$$v = r \frac{d\phi}{dt}$$

$$p = \sum_n p_n = \left(\sum_n \rho_n R_n \right) T = \rho \left(\sum_n \frac{\rho_n R_n}{\rho} \right) T = \rho \left(\sum_n q_n R_n \right) T = \rho RT$$

Gas law

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho u}{r \cos \phi \partial \lambda} + \frac{\partial \rho v \cos \phi}{r \cos \phi \partial \phi} + \frac{\partial \rho r^2 w}{r^2 \partial r} = F_\rho$$

Density

From IEE

$$\frac{\partial \rho e}{\partial t} + \nabla \cdot \rho e V + p \nabla \cdot V = \rho Q$$

replace

$$e = \sum_{i=1}^N q_i e_i = \sum_{i=1}^N q_i C_{V_i} T = C_V T = (C_P - R) T = h - RT$$

We have

$$\frac{\partial \rho h}{\partial t} + \nabla \cdot \rho h V - \frac{dp}{dt} = \rho Q$$

Specific value

where $C_V = \sum_{i=1}^N q_i C_{V_i}$ $C_P = \sum_{i=1}^N q_i C_{P_i}$ $R = \sum_{i=1}^N q_i R_i$ & $q_i = \frac{\rho_i}{\rho}$ $\rho = \sum_{n=1}^N \rho_n$

Combine IEE in h form with

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho V = F_\rho$$

We have

$$\frac{dh}{dt} - \frac{\kappa h}{p} \frac{dp}{dt} = Q - \frac{h}{\rho} F_\rho$$

& we need

$$\frac{dq_i}{dt} = F_{q_i}$$

where $h=CpT$ is an Enthalpy

$$\frac{du}{dt} - \frac{uv \tan \phi}{r} + \frac{uw}{r} - f_s v + f_c w + \frac{\kappa h}{p} \frac{1}{r \cos \phi} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \lambda} \right) = F_u$$

$$\frac{dv}{dt} + \frac{u^2 \tan \phi}{r} + \frac{vw}{r} + f_s u + \frac{\kappa h}{p} \frac{1}{r} \left(\frac{\partial p}{\partial \phi} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \phi} \right) = F_v$$

$$\frac{dw}{dt} - \frac{u^2 + v^2}{r} - f_c u + \frac{\kappa h}{p} \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} + g = F_w$$

$$\frac{dh}{dt} - \frac{\kappa h}{p} \frac{dp}{dt} = F_h$$

$$\frac{\partial \beta}{\partial t} + \frac{\partial \beta}{\partial \lambda} \dot{\lambda} + \frac{\partial \beta}{\partial \phi} \dot{\phi} + \frac{\partial \beta}{\partial \zeta} \dot{\zeta} = F_\rho^\beta$$

$$\frac{dq_i}{dt} = F_{q_i}$$

Deep Atmos Dyn in generalized coordinate

Staniforth and Wood (2003)

where $\frac{d()}{dt} = \frac{\partial ()}{\partial t} + \dot{\lambda} \frac{\partial ()}{\partial \lambda} + \dot{\phi} \frac{\partial ()}{\partial \phi} + \dot{\zeta} \frac{\partial ()}{\partial \zeta}$; $\beta = \rho r^2 \cos \phi \frac{\partial r}{\partial \zeta}$; $w = \frac{\partial r}{\partial t} + \dot{\lambda} \frac{\partial r}{\partial \lambda} + \dot{\phi} \frac{\partial r}{\partial \phi} + \dot{\zeta} \frac{\partial r}{\partial \zeta}$
 $f_s = 2\Omega \sin \phi$; $f_c = 2\Omega \cos \phi$; $g = g(r)$; $p = \rho \kappa h$

Angular momentum

$$A = r \cos \phi (u + \Omega r \cos \phi)$$

$$\frac{dA}{dt} = r \cos \phi \frac{du}{dt} + (u + 2\Omega r \cos \phi) \frac{dr \cos \phi}{dt} = r \cos \phi \frac{du}{dt} + (u + 2\Omega r \cos \phi) (w \cos \phi - v \sin \phi)$$

put $\frac{du}{dt} = \frac{uv \tan \phi}{r} - \frac{uw}{r} + f_s v - f_c w - \frac{\kappa h}{p} \frac{1}{r \cos \phi} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} \frac{\partial r}{\partial \lambda} \right) + F_u$ in

We have $\frac{dA}{dt} = r \cos \phi \left[F_u - \frac{\kappa h}{p} \frac{1}{r \cos \phi} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} \frac{\partial r}{\partial \lambda} \right) \right]$

$$\frac{dA}{dt} = \frac{\partial A}{\partial t} + \frac{u}{r \cos \phi} \frac{\partial A}{\partial \lambda} + \frac{v}{r} \frac{\partial A}{\partial \phi} + \dot{\xi} \frac{\partial A}{\partial \xi}$$

$$\frac{\partial \beta}{\partial t} + \frac{\partial \frac{u}{r \cos \phi} \beta}{\partial \lambda} + \frac{\partial \frac{v}{r} \beta}{\partial \phi} + \frac{\partial \dot{\xi} \beta}{\partial \xi} = 0$$

$$\frac{\partial \beta A}{\partial t} + \frac{\partial \frac{u}{r \cos \phi} \beta A}{\partial \lambda} + \frac{\partial \frac{v}{r} \beta A}{\partial \phi} + \frac{\partial \dot{\xi} \beta A}{\partial \xi} = \beta \frac{dA}{dt}$$

where

$$\beta = \rho r^2 \cos \phi \frac{\partial r}{\partial \xi}$$

$$\begin{aligned}
& \frac{\partial \beta A}{\partial t} + \frac{\partial}{\partial \lambda} \left(\frac{u \beta A}{r \cos \phi} \right) + \frac{\partial}{\partial \phi} \left(\frac{v \beta A}{r} \right) + \frac{\partial \dot{\xi} \beta A}{\partial \xi} \\
&= \beta r \cos \phi \left[F_u - \frac{\kappa h}{p} \frac{1}{r \cos \phi} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} \frac{\partial r}{\partial \lambda} \right) \right] \\
&= -r^2 \cos \phi \frac{\partial r}{\partial \xi} \frac{\partial p}{\partial \lambda} + r^2 \cos \phi \frac{\partial p}{\partial \xi} \frac{\partial r}{\partial \lambda} + \beta r \cos \phi F_u \\
&= -\frac{\partial p r^2 \cos \phi \frac{\partial r}{\partial \xi}}{\partial \lambda} + p \frac{\partial r^2 \cos \phi \frac{\partial r}{\partial \xi}}{\partial \lambda} + \frac{\partial p r^2 \cos \phi \frac{\partial r}{\partial \lambda}}{\partial \xi} - p \frac{\partial r^2 \cos \phi \frac{\partial r}{\partial \lambda}}{\partial \xi} + \beta r \cos \phi F_u \\
&= -\frac{\partial p r^2 \cos \phi \frac{\partial r}{\partial \xi}}{\partial \lambda} + \frac{\partial p r^2 \cos \phi \frac{\partial r}{\partial \lambda}}{\partial \xi} + \beta r \cos \phi F_u
\end{aligned}$$

$$\frac{\partial}{\partial t} \iiint \rho A dv = \iiint \rho r \cos \phi F_u dv + \iint \left(p \frac{\partial r}{\partial \lambda} \right)_T ds - \iint \left(p \frac{\partial r}{\partial \lambda} \right)_B ds$$

Angular momentum conserved if top $p=0$ or $r=\text{constant}$
thus the bottom is the only torque.

Consider Kinetic energy from momentum equations

$$\begin{aligned}
 u \frac{du}{dt} - u \frac{uv \tan \phi}{r} + u \frac{uw}{r} - uf_s v + uf_c w + u \frac{\kappa h}{p} \frac{1}{r \cos \phi} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \lambda} \right) &= 0 \\
 v \frac{dv}{dt} + v \frac{u^2 \tan \phi}{r} + v \frac{vw}{r} + vf_s u + v \frac{\kappa h}{p} \frac{1}{r} \left(\frac{\partial p}{\partial \phi} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \phi} \right) &= 0 \\
 w \frac{dw}{dt} - w \frac{u^2 + v^2}{r} - wf_c u + w \frac{\kappa h}{p} \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} + wg &= 0
 \end{aligned}$$

Sum them together, we have $K = \frac{1}{2}(u^2 + v^2 + w^2)$

$$\frac{\partial K}{\partial t} + \dot{\lambda} \frac{\partial K}{\partial \lambda} + \dot{\phi} \frac{\partial K}{\partial \phi} + \dot{\zeta} \frac{\partial K}{\partial \zeta} = -\frac{1}{\rho} \dot{\lambda} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \lambda} \right) - \frac{1}{\rho} \dot{\phi} \left(\frac{\partial p}{\partial \phi} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \phi} \right) - \frac{1}{\rho} w \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} - gw$$

Combine with mass and w equations as

$$\frac{\partial \beta}{\partial t} + \frac{\partial \beta \dot{\lambda}}{\partial \lambda} + \frac{\partial \beta \dot{\phi}}{\partial \phi} + \frac{\partial \beta \dot{\zeta}}{\partial \zeta} = 0 \quad \text{and} \quad w = \frac{\partial r}{\partial t} + \dot{\lambda} \frac{\partial r}{\partial \lambda} + \dot{\phi} \frac{\partial r}{\partial \phi} + \dot{\zeta} \frac{\partial r}{\partial \zeta} \quad \text{We have}$$

$$\begin{aligned}
 \frac{\partial \beta K}{\partial t} + \frac{\partial \dot{\lambda} \beta \left(K + \frac{p}{\rho} \right)}{\partial \lambda} + \frac{\partial \dot{\phi} \beta \left(K + \frac{p}{\rho} \right)}{\partial \phi} + \frac{\partial \dot{\zeta} \beta \left(K + \frac{p}{\rho} \right)}{\partial \zeta} &= p \frac{\partial \beta \dot{\lambda}}{\partial \lambda} + p \frac{\partial \beta \dot{\phi}}{\partial \phi} + p \frac{\partial \beta \dot{\zeta}}{\partial \zeta} - \frac{\partial p}{\rho} \frac{\partial r}{\partial t} \frac{\partial \zeta}{\partial r} + p \frac{\partial \beta}{\rho} \frac{\partial r}{\partial t} \frac{\partial \zeta}{\partial r} - \beta gw \\
 &= p \left(\frac{\partial \beta \dot{\lambda}}{\partial \lambda} + \frac{\partial \beta \dot{\phi}}{\partial \phi} + \frac{\partial \beta \dot{\zeta}}{\partial \zeta} \right) - \frac{\partial p r^2 \cos \phi}{\partial \zeta} \frac{\partial r}{\partial t} + p \frac{\partial r^2 \cos \phi}{\partial \zeta} \frac{\partial r}{\partial t} - \beta gw = p \left(\frac{\partial \beta}{\partial t} + \frac{\partial \beta \dot{\lambda}}{\partial \lambda} + \frac{\partial \beta \dot{\phi}}{\partial \phi} + \frac{\partial \beta \dot{\zeta}}{\partial \zeta} \right) - \frac{\partial p r^2 \cos \phi}{\partial \zeta} \frac{\partial r}{\partial t} - \beta gw
 \end{aligned}$$

Consider geo-potential energy

$$\int g dr = \Phi \quad \text{or} \quad g = \frac{d\Phi}{dr}$$

$$g_w = \frac{d\Phi}{dr} \frac{dr}{dt} = \frac{\partial\Phi}{\partial t} + \dot{\lambda} \frac{\partial\Phi}{\partial\lambda} + \dot{\phi} \frac{\partial\Phi}{\partial\phi} + \dot{\xi} \frac{\partial\Phi}{\partial\xi}$$

Combine with mass equation $\frac{\partial\beta}{\partial t} + \frac{\partial\beta\dot{\lambda}}{\partial\lambda} + \frac{\partial\beta\dot{\phi}}{\partial\phi} + \frac{\partial\beta\dot{\xi}}{\partial\xi} = 0$

We have $\frac{\partial\beta\Phi}{\partial t} + \frac{\partial\dot{\lambda}\beta\Phi}{\partial\lambda} + \frac{\partial\dot{\phi}\beta\Phi}{\partial\phi} + \frac{\partial\dot{\xi}\beta\Phi}{\partial\xi} = \beta g_w$

Consider internal energy

$$\frac{dC_p T}{dt} - \frac{1}{\rho} \frac{dp}{dt} = \frac{dC_v T}{dt} + \frac{dT}{dt} - \frac{1}{\rho} \frac{dp}{dt} = \frac{dC_v T}{dt} + p \frac{d\frac{1}{\rho}}{dt} = 0$$

Apply mass equations in both total derivative terms,
We have

$$\frac{\partial\beta C_v T}{\partial t} + \frac{\partial\dot{\lambda}\beta C_v T}{\partial\lambda} + \frac{\partial\dot{\phi}\beta C_v T}{\partial\phi} + \frac{\partial\dot{\xi}\beta C_v T}{\partial\xi} = -p \left[\frac{\partial}{\partial t} \left(\frac{\beta}{\rho} \right) + \frac{\partial}{\partial\lambda} \left(\dot{\lambda} \frac{\beta}{\rho} \right) + \frac{\partial}{\partial\phi} \left(\dot{\phi} \frac{\beta}{\rho} \right) + \frac{\partial}{\partial\xi} \left(\dot{\xi} \frac{\beta}{\rho} \right) \right]$$

Combine previous three red enclosed equations

$$\frac{\partial \beta K}{\partial t} + \frac{\partial \dot{\lambda} \beta \left(K + \frac{P}{\rho} \right)}{\partial \lambda} + \frac{\partial \dot{\phi} \beta \left(K + \frac{P}{\rho} \right)}{\partial \phi} + \frac{\partial \dot{\xi} \beta \left(K + \frac{P}{\rho} \right)}{\partial r} = p \left(\frac{\partial \frac{\beta}{\rho}}{\partial t} + \frac{\partial \frac{\beta}{\rho} \dot{\lambda}}{\partial \lambda} + \frac{\partial \frac{\beta}{\rho} \dot{\phi}}{\partial \phi} + \frac{\partial \frac{\beta}{\rho} \dot{\xi}}{\partial r} \right) - \frac{\partial p r^2 \cos \phi \frac{\partial r}{\partial t}}{\partial \xi} - \beta g w$$

$$\frac{\partial \beta \Phi}{\partial t} + \frac{\partial \dot{\lambda} \beta \Phi}{\partial \lambda} + \frac{\partial \dot{\phi} \beta \Phi}{\partial \phi} + \frac{\partial \dot{\xi} \beta \Phi}{\partial r} = \beta g w$$

$$\frac{\partial \beta C_v T}{\partial t} + \frac{\partial \dot{\lambda} \beta C_v T}{\partial \lambda} + \frac{\partial \dot{\phi} \beta C_v T}{\partial \phi} + \frac{\partial \dot{\xi} \beta C_v T}{\partial r} = -p \left[\frac{\partial}{\partial t} \left(\frac{\beta}{\rho} \right) + \frac{\partial}{\partial \lambda} \left(\dot{\lambda} \frac{\beta}{\rho} \right) + \frac{\partial}{\partial \phi} \left(\dot{\phi} \frac{\beta}{\rho} \right) + \frac{\partial}{\partial r} \left(\dot{\xi} \frac{\beta}{\rho} \right) \right]$$

Integral globally with all BC, include $p_{\xi_T} = 0$ $\left(\frac{\partial r}{\partial t} \right)_{\xi_B} = 0$

$$\iiint \left(\frac{\partial \beta K}{\partial t} + \frac{\partial \beta C_v T}{\partial t} + \frac{\partial \beta \Phi}{\partial t} \right) d\xi d\lambda d\phi = - \iint \left[\left(p r^2 \cos \phi \frac{\partial r}{\partial t} \right)_{\xi_T} - \left(p r^2 \cos \phi \frac{\partial r}{\partial t} \right)_{\xi_B} \right] d\lambda d\phi$$

and BC give $\frac{\partial}{\partial t} \iiint \beta (K + C_v T + \Phi) d\xi d\lambda d\phi = \frac{\partial}{\partial t} \iiint \rho (K + C_v T + \Phi) dv = 0$

Total energy conserved

Deep Atmos vs non-Hydro

- From Deep atmosphere, we require r changes with time, thus we need dw/dt equation
- And we need full curvature and Coriolis force terms to satisfy conservation
- Thus, based on conservation requirement, **a deep atmospheric dynamic is a non-hydrostatic dynamic. A non-hydrostatic dynamics can be shallow or deep atmospheric dynamics.**
- Both r and vertical components of curvature and Coriolis force should be considered in deep atmosphere; and should not be considered in shallow atmosphere.

$$\frac{du^*}{dt} + \frac{u^* w}{r} - f_s v^* + f_c^* w + \frac{\kappa h}{p} \frac{1}{r} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \lambda} \right) = F_u$$

$$\frac{dv^*}{dt} + \frac{v^* w}{r} + f_s u^* + m^2 \frac{s^{*2}}{r} \sin \phi + \frac{\kappa h}{p} \frac{1}{r} \left(\frac{\partial p}{\partial \varphi} - \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} \frac{\partial r}{\partial \varphi} \right) = F_v$$

$$\frac{dw}{dt} - m^2 \frac{s^{*2}}{r} - m^2 f_c^* u^* + \frac{\kappa h}{p} \frac{\partial p}{\partial \zeta} \frac{\partial \zeta}{\partial r} + g = F_w$$

$$\frac{dh}{dt} - \frac{\kappa h}{p} \frac{dp}{dt} = F_h$$

Deep Atmos Dyn
in spherical mapping
& generalized coordinates

$$\frac{\partial \rho^*}{\partial t} + m^2 \frac{\partial \rho^*}{\partial \lambda} \frac{u^*}{r} + m^2 \frac{\partial \rho^*}{\partial \varphi} \frac{v^*}{r} + \frac{\partial \rho^*}{\partial \zeta} \dot{\zeta} = F_\rho^*$$

Staniforth and Wood (2003)
Juang (2014) NCEP Office Note

$$\frac{dq_i}{dt} = F_{q_i}$$

$$p = \rho \kappa h$$

where $\frac{d()}{dt} = \frac{\partial ()}{\partial t} + \dot{\lambda} \frac{\partial ()}{\partial \lambda} + \dot{\varphi} \frac{\partial ()}{\partial \varphi} + \dot{\zeta} \frac{\partial ()}{\partial \zeta} = \frac{\partial ()}{\partial t} + m^2 u^* \frac{\partial ()}{r \partial \lambda} + m^2 v^* \frac{\partial ()}{r \partial \varphi} + \dot{\zeta} \frac{\partial ()}{\partial \zeta}$; $\rho^* = \rho \frac{r^2}{a^2} \frac{\partial r}{\partial \zeta}$

$$f_s = 2\Omega \sin \phi ; f_c^* = 2\Omega \cos^2 \phi ; g = g(r) ; \kappa = \frac{R}{C_p} ; \gamma = \frac{C_p}{C_v} ; s^{*2} = u^{*2} + v^{*2}$$

Start from continuity equation to have similar to shallow and hydrostatic system to make as mass coordinates, Mass at give area can be obtained by integral vertical as

$$Mass = \int_{\phi_1}^{\phi_2} \int_{\lambda_1}^{\lambda_2} \int_{\xi_{sfc}}^{\xi_{Top}} \rho r^2 \cos \phi \frac{\partial r}{\partial \xi} d\xi d\lambda d\phi$$

Then project this mass to earth surface

$$pressure = \frac{Mass \bar{g}}{Area} = \frac{\iiint \rho \bar{g} r^2 \cos \phi \frac{\partial r}{\partial \xi} d\xi d\lambda d\phi}{\iint a^2 \cos \phi d\lambda d\phi}$$

We define it as coordinate pressure

Previous integral can be deduced to be only in vertical as

$$\int_{\xi_{surface}}^{\xi_{TOP}} \rho \bar{g} \frac{r^2}{a^2} \frac{\partial r}{\partial \xi} d\xi = \tilde{p}_{\xi_{surface}} = - \int_{\xi_{surface}}^{\xi_{TOP}} \frac{\partial \tilde{p}}{\partial \xi} d\xi$$

Thus, we have
$$\frac{\partial \tilde{p}}{\partial \xi} = -\rho \bar{g} \frac{r^2}{a^2} \frac{\partial r}{\partial \xi} = -\rho^* \bar{g}$$

Put into continuity equation

$$\frac{\partial \rho^*}{\partial t} + m^2 \frac{\partial \rho^* \frac{u^*}{r}}{\partial \lambda} + m^2 \frac{\partial \rho^* \frac{v^*}{r}}{\partial \varphi} + \frac{\partial \rho^* \dot{\xi}}{\partial \xi} = 0$$

We have

$$\frac{\partial \frac{\partial \tilde{p}}{\partial \xi}}{\partial t} + m^2 \left(\frac{\partial \frac{\partial \tilde{p}}{\partial \xi} \frac{u^*}{r}}{\partial \lambda} + \frac{\partial \frac{\partial \tilde{p}}{\partial \xi} \frac{v^*}{r}}{\partial \varphi} \right) + \frac{\partial \frac{\partial \tilde{p}}{\partial \xi} \dot{\xi}}{\partial \xi} = 0$$

Since $\frac{\partial \tilde{p}}{\partial \xi} = -\rho^* g$ and $\rho^* > 0$

Thus \tilde{p} is monotone with vertical coordinate

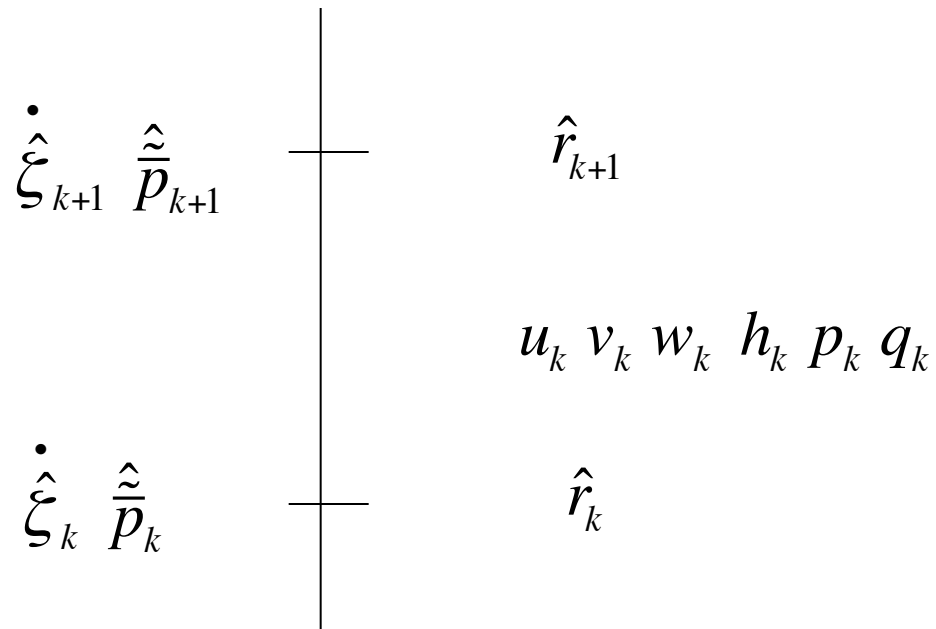
We can use it for coordinate definition and we can call it height weighted coordinate pressure or simply coordinate pressure

So we can have $\hat{\tilde{p}}_k = \hat{A}_k + \hat{B}_k \tilde{p}_s + \hat{C}_k \left(\frac{h_{k-1} + h_k}{h_{0k-1} + h_{0k}} \right)^{C_{pd}/R_d}$

For opr compatibility, we use

$$\hat{\tilde{p}}_k = \hat{A}_k + \hat{B}_k \tilde{p}_s$$

$k=K$, last layer at the top



$k=1$, the first layer next to ground

For dynamics, all source terms are set to be zero !

Vertical integral angular momentum principle, we have

$$\begin{aligned}
 \int_{\xi_s}^{\xi_T} \frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \left(\frac{\partial p}{\partial \lambda} - \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} \frac{\partial r}{\partial \lambda} \right) d\xi &= \int_{\xi_s}^{\xi_T} \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial p}{\partial \lambda} - \frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} \frac{\partial r}{\partial \lambda} \right) d\xi \\
 &= \int_{\xi_s}^{\xi_T} \left(-\frac{\partial r}{\partial \xi} \frac{\bar{g} r^2}{a^2} \frac{\partial p}{\partial \lambda} + \frac{\bar{g} r^2}{a^2} \frac{\partial p}{\partial \xi} \frac{\partial r}{\partial \lambda} \right) d\xi \\
 &= \int_{\xi_s}^{\xi_T} \frac{\bar{g}}{3a^2} \left(-\frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} + \frac{\partial p}{\partial \xi} \frac{\partial r^3}{\partial \lambda} \right) d\xi \\
 &= \int_{\xi_s}^{\xi_T} \frac{\bar{g}}{3a^2} \left(-\frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} + \frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} - r^3 \frac{\partial}{\partial \xi} \frac{\partial p}{\partial \lambda} \right) d\xi \\
 &= \int_{\xi_s}^{\xi_T} \frac{\bar{g}}{3a^2} \left(\frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} - \frac{\partial r^3}{\partial \lambda} \frac{\partial p}{\partial \xi} \right) d\xi \\
 &= \frac{\partial}{\partial \lambda} \int_{\xi_s}^{\xi_T} \frac{\bar{g}}{3a^2} r^3 \frac{\partial p}{\partial \xi} d\xi - \frac{\bar{g} r_T^3}{3a^2} \frac{\partial p_T}{\partial \lambda} + \frac{\bar{g} r_S^3}{3a^2} \frac{\partial p_S}{\partial \lambda}
 \end{aligned}$$

Discretize the fourth and last LHSs

$$\int_{\xi_S}^{\xi_T} \frac{\bar{g}}{3a^2} \left(-\frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} + \frac{\partial r^3}{\partial \xi} \frac{\partial p}{\partial \lambda} - r^3 \frac{\partial}{\partial \xi} \frac{\partial p}{\partial \lambda} \right) d\xi = \frac{\partial}{\partial \lambda} \int_{\xi_S}^{\xi_T} \frac{\bar{g}}{3a^2} r^3 \frac{\partial p}{\partial \xi} d\xi - \frac{\bar{g}r_T^3}{3a^2} \frac{\partial p_T}{\partial \lambda} + \frac{\bar{g}r_S^3}{3a^2} \frac{\partial p_S}{\partial \lambda}$$

With P top zero, we have

$$\sum_{k=1}^K \left[(\hat{r}_{k+1}^3 - \hat{r}_k^3) \frac{\partial p_k}{\partial \lambda} + r_k^3 \left(\frac{\partial \hat{p}_{k+1}}{\partial \lambda} - \frac{\partial \hat{p}_k}{\partial \lambda} \right) \right] = -\hat{r}_S^3 \frac{\partial p_S}{\partial \lambda}$$

let $p_k = f(\hat{p}_{k+1}, \hat{p}_k)$ so $p_k = \frac{\partial f_k}{\partial \hat{p}_{k+1}} \hat{p}_{k+1} + \frac{\partial f_k}{\partial \hat{p}_k} \hat{p}_k$ and $\frac{\partial f_k}{\partial \hat{p}_{k+1}} + \frac{\partial f_k}{\partial \hat{p}_k} = 1$

Then we have

$$\sum_{k=1}^K \left[\left\langle (\hat{r}_{k+1}^3 - \hat{r}_k^3) \frac{\partial f_k}{\partial \hat{p}_{k+1}} + r_k^3 \right\rangle \frac{\partial \hat{p}_{k+1}}{\partial \lambda} + \left\langle (\hat{r}_{k+1}^3 - \hat{r}_k^3) \frac{\partial f_k}{\partial \hat{p}_k} - r_k^3 \right\rangle \frac{\partial \hat{p}_k}{\partial \lambda} \right] = -\hat{r}_S^3 \frac{\partial p_S}{\partial \lambda}$$

For simplicity, we let

$$\frac{\partial f_k}{\partial \hat{p}_k} = \frac{\partial f_k}{\partial \hat{p}_{k+1}} = \frac{1}{2} \quad \text{SO} \quad r_k^3 = \frac{1}{2} \hat{r}_{k+1}^3 + \frac{1}{2} \hat{r}_k^3$$

$$\text{since} \quad \hat{r}_{k+1}^3 = \hat{r}_k^3 + 3 \left(\frac{\kappa h}{p \bar{g}} \right)_k a^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right)$$

$$\text{SO} \quad r_k^3 = \hat{r}_k^3 + \frac{3}{2} \left(\frac{\kappa h}{p \bar{g}} \right)_k a^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right)$$

remove layer value we have

$$\hat{r}_{k+1}^3 = \hat{r}_k^3 + 3 \frac{a^2}{\bar{g}} \kappa_k h_k \frac{\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1}}{p_k}$$

SO

$$r_k^3 = \hat{r}_1^3 + 3 \frac{a^2}{\bar{g}} \sum_{i=1}^{k-1} \left[\kappa_i h_i \frac{\hat{\bar{p}}_i - \hat{\bar{p}}_{i+1}}{p_i} \right] + \frac{3}{2} \frac{a^2}{\bar{g}} \kappa_k h_k \frac{\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1}}{p_k}$$

We have

$$\boxed{\frac{\hat{r}_{k+1}^3}{a^2} = \frac{\hat{r}_k^3}{a^2} + 3 \left(\frac{\kappa h}{p \bar{g}} \right)_k \left(\hat{p}_k - \hat{p}_{k+1} \right)}$$

then

transform to spectral for r to do derivative, then
transform back to grid space for nonlinear computing.

$$\frac{du_k^*}{dt} = -\frac{u_k^* w_k}{r_k} + f_s v_k^* - f_c^* w_k - \frac{\kappa_k h_k}{p_k} \frac{a}{r_k} \frac{\partial p_k}{a \partial \lambda} - \frac{\bar{g}}{6} \frac{\hat{p}_k - \hat{p}_{k+1}}{\hat{p}_k - \hat{p}_{k+1}} \frac{a}{r_k} \frac{\partial \left(\frac{\hat{r}_k^3}{a^2} + \frac{\hat{r}_{k+1}^3}{a^2} \right)}{a \partial \lambda}$$

$$\frac{dv_k^*}{dt} = -\frac{v_k^* w_k}{r_k} - f_s u_k^* - m^2 \frac{s_k^{*2}}{r_k} \sin \phi - \frac{\kappa_k h_k}{p_k} \frac{a}{r_k} \frac{\partial p_k}{a \partial \varphi} - \frac{\bar{g}}{6} \frac{\hat{p}_k - \hat{p}_{k+1}}{\hat{p}_k - \hat{p}_{k+1}} \frac{a}{r_k} \frac{\partial \left(\frac{\hat{r}_k^3}{a^2} + \frac{\hat{r}_{k+1}^3}{a^2} \right)}{a \partial \varphi}$$

Select two forms of advection of PGF

$$\begin{aligned} \frac{1}{\gamma} \frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{dp}{dt} &= -p \nabla_H \cdot \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} V_H \right) + \frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial p}{\partial \xi} \frac{\partial \zeta}{\partial r} (w - V_H \cdot \nabla_H r) - \frac{\partial}{\partial \xi} \left(p \frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial \zeta}{\partial r} (w - V_H \cdot \nabla_H r) \right) \\ &= -p \nabla_H \cdot \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} V_H \right) - p \frac{\partial}{\partial \xi} \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial \zeta}{\partial r} (w - V_H \cdot \nabla_H r) \right) \end{aligned}$$

Continue discretization, we have

$$\frac{dp}{dt} = -\gamma p \frac{1}{\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p}} \nabla_H \cdot \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} V_H \right) - \gamma p \frac{1}{\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p}} \frac{\partial}{\partial \xi} \left(\frac{\partial \tilde{p}}{\partial \xi} \frac{\kappa h}{p} \frac{\partial \zeta}{\partial r} (w - V_H \cdot \nabla_H r) \right)$$

$$\begin{aligned} \left(\frac{dp}{dt} \right)_k &= -\gamma_k p_k \left(\nabla_H \cdot (V_H) + \frac{1}{\Delta r^3} V_H \cdot \nabla_H \Delta r^3 - \frac{1}{\Delta r^3} \Delta (V_H \cdot \nabla_H r^3) + \frac{a^2}{r^2} \frac{\Delta \tilde{w}}{\Delta r} \right)_k \\ &= -\gamma_k p_k \left(\left(\frac{\partial \dot{\lambda}}{\partial \lambda} + \frac{\partial \dot{\mu}}{\partial \mu} \right)_k - \frac{\frac{\dot{\lambda}_{k-1} - \dot{\lambda}_k}{2} \frac{\partial \hat{r}_k^3}{\partial \lambda} + \frac{\dot{\lambda}_k - \dot{\lambda}_{k+1}}{2} \frac{\partial \hat{r}_{k+1}^3}{\partial \lambda} + \frac{\dot{\mu}_{k-1} - \dot{\mu}_k}{2} \frac{\partial \hat{r}_k^3}{\partial \mu} + \frac{\dot{\mu}_k - \dot{\mu}_{k+1}}{2} \frac{\partial \hat{r}_{k+1}^3}{\partial \mu}}{\Delta r^3} + \left(\frac{a^2}{r^2} \frac{\Delta \tilde{w}}{\Delta r} \right)_k \right) \\ &= -\gamma_k p_k \left(m^2 \left[\left(\frac{\partial \frac{u^*}{r}}{\partial \lambda} + \frac{\partial \frac{v^*}{r}}{\partial \varphi} \right)_k - \frac{\left(\frac{u_{k-1}^*}{r_{k-1}} - \frac{u_k^*}{r_k} \right) \frac{\partial \hat{r}_k^3}{\partial \lambda} + \left(\frac{u_k^*}{r_k} - \frac{u_{k+1}^*}{r_{k+1}} \right) \frac{\partial \hat{r}_{k+1}^3}{\partial \lambda} + \left(\frac{v_{k-1}^*}{r_{k-1}} - \frac{v_k^*}{r_k} \right) \frac{\partial \hat{r}_k^3}{\partial \varphi} + \left(\frac{v_k^*}{r_k} - \frac{v_{k+1}^*}{r_{k+1}} \right) \frac{\partial \hat{r}_{k+1}^3}{\partial \varphi}}{2 \Delta r^3} \right] + \left(\frac{a^2}{r^2} \frac{\Delta \tilde{w}}{\Delta r} \right)_k \right) \end{aligned}$$

So we use dp/dt from momentum equation to thermodynamic equation, thus total energy will be conserved.

$$\frac{dh}{dt} - \frac{\kappa h}{p} \frac{dp}{dt} = 0$$

SO

$$\frac{dh_k}{dt} = -\gamma_k \kappa_k h_k \left\{ m^2 \left(\frac{\partial \frac{u^*}{r}}{\partial \lambda} + \frac{\partial \frac{v^*}{r}}{\partial \varphi} \right)_k - m^2 \frac{\left(\frac{u_{k-1}^*}{r_{k-1}} - \frac{u_k^*}{r_k} \right) \frac{\partial \hat{r}_k^3}{\partial \lambda} + \left(\frac{u_k^*}{r_k} - \frac{u_{k+1}^*}{r_{k+1}} \right) \frac{\partial \hat{r}_{k+1}^3}{\partial \lambda} + \left(\frac{v_{k-1}^*}{r_{k-1}} - \frac{v_k^*}{r_k} \right) \frac{\partial \hat{r}_k^3}{\partial \varphi} + \left(\frac{v_k^*}{r_k} - \frac{v_{k+1}^*}{r_{k+1}} \right) \frac{\partial \hat{r}_{k+1}^3}{\partial \varphi}}{2\Delta r^3} + \left(\frac{a^2}{r^2} \frac{\Delta \tilde{w}}{\Delta r} \right)_k \right\}$$

where $\tilde{w} = \frac{r^2}{a^2} w$

$$\frac{dw}{dt} = m^2 \frac{s^{*2}}{r} + m^2 f_c^* u^* - \frac{\kappa h}{p} \frac{\partial p}{\partial \xi} \frac{\partial \xi}{\partial r} - g = m^2 \frac{s^{*2}}{r} + m^2 f_c^* u^* + \bar{g} \frac{r^2 \Delta p}{a^2 \Delta \tilde{p}} - g$$

$$\frac{dr^2 w}{dt} = 2rw^2 + r^2 \frac{dw}{dt} = 2 \frac{(r^2 w)^2}{r^3} + m^2 r s^{*2} + m^2 r^2 f_c^* u^* + \bar{g} a^2 \left(\frac{r^4 \Delta p}{a^4 \Delta \tilde{p}} - 1 \right)$$

and $\tilde{w} = \frac{r^2}{a^2} w$

so
$$\frac{d\tilde{w}}{dt} = 2a^2 \frac{\tilde{w}^2}{r^3} + m^2 \frac{r}{a^2} s^{*2} + m^2 \frac{r^2}{a^2} f_c^* u^* + \bar{g} \left(\frac{r^4 \Delta p}{a^4 \Delta \tilde{p}} - 1 \right)$$

Since $r_k^2 w_k = \frac{1}{2} \hat{r}_{k+1}^2 \hat{w}_{k+1} + \frac{1}{2} \hat{r}_k^2 \hat{w}_k$ so $\tilde{w}_k = \frac{1}{2} \hat{w}_{k+1} + \frac{1}{2} \hat{w}_k$

And BC
$$\hat{w}_1 = m^2 \frac{u_1^*}{\hat{r}_1} \frac{\partial \hat{r}_1}{\partial \lambda} + m^2 \frac{v_1^*}{\hat{r}_1} \frac{\partial \hat{r}_1}{\partial \varphi}$$

so
$$\hat{w}_1 = m^2 \frac{\hat{r}_1}{a} \left(u_1^* \frac{\partial \hat{r}_1}{a \partial \lambda} + v_1^* \frac{\partial \hat{r}_1}{a \partial \varphi} \right)$$
 Then we have w at all levels

Linearize for SISL

- Collect discretized equations
- Linearization
- Matrixes for semi-implicit
- Add semi-Lagrangian with semi-implicit

To linearize, we define a base state

start

$$\hat{p}_{01} = 101.326 \quad ; \quad \hat{r}_{01} = a = 6371220 \quad ; \quad h_{0k} = C_{Pd} T_{0k} \quad ; \quad T_{0k} = 300$$

$$\hat{\tilde{p}}_{0k} = \hat{A}_k + \hat{B}_k \hat{p}_{01} \quad ; \quad \Delta \tilde{p}_{0k} = \hat{\tilde{p}}_{0k} - \hat{\tilde{p}}_{0k+1} \quad ; \quad k_0 = R_d / C_{Pd}$$

Since p equals to p bar, so
$$\frac{\partial \tilde{p}}{\partial \xi} = -\rho g \frac{r^2}{a^2} \frac{\partial r}{\partial \xi} = -\frac{\tilde{p} g}{\kappa h} \frac{\partial r}{\partial \xi}$$

So we can get all r at interface as
$$\hat{r}_{0k+1} = \hat{r}_{0k} + \frac{\kappa_{0k} h_{0k}}{\bar{g}} \ln \frac{\hat{\tilde{p}}_{0k}}{\hat{\tilde{p}}_{0k+1}}$$

then use
$$\hat{p}_{0k} - \hat{p}_{0k+1} = \left(\hat{\tilde{p}}_{0k} - \hat{\tilde{p}}_{0k+1} \right) \frac{a^2}{r_{0k}^2} \quad ; \quad p_{0k} = \frac{1}{2} \left(\hat{p}_{0k} + \hat{p}_{0k+1} \right)$$

then
$$r_{0k}^3 = \frac{1}{2} \left(\hat{r}_{0k}^3 + \hat{r}_{0k+1}^3 \right) \quad \Delta p_{0k} = \hat{p}_{0k} - \hat{p}_{0k+1}$$

and
$$\varepsilon_{0k} = \frac{r_{0k}}{a}$$

Summary of linear equations in matrix/vector short form

$$\left(\frac{D\tilde{p}_s}{Dt} \right)_L = -\Pi_{1i} D_i^*$$

$$\left(\frac{d\tilde{w}_k}{dt} \right)_L = -b_{0k} \tilde{p}_s + \Gamma_{ki} p_i$$

$$\left(\frac{dp_k}{dt} \right)_L = -d_{0k} D_k^* + M_{ki} \tilde{w}_i$$

$$\left(\frac{dh_k}{dt} \right)_L = -f_{0k} D_k^* + Z_{ki} \tilde{w}_i$$

$$\left(\frac{dD_k^*}{dt} \right)_L = \frac{n(n+1)}{a^2} (A_{ki} p_i + B_{ki} h_i + e_k p_s)$$

$$\delta D_k^* - \frac{n(n+1)}{a^2} \alpha \delta t \left(A_{ki} \delta p_i + B_{ki} \delta h_i + e_k \delta \tilde{p}_s \right) = S_{D^*}$$

$$\delta \tilde{w}_k + \alpha \delta t \left(b_{0k} \delta \tilde{p}_s - \Gamma_{ki} \delta p_i \right) = S_{\tilde{w}_k}$$

$$\delta h_k + \alpha \delta t \left(f_{0k} \delta D_k^* - Z_{ki} \delta \tilde{w}_i \right) = S_{h_k}$$

$$\delta p_k + \alpha \delta t \left(d_{0k} \delta D_k^* - M_{ki} \delta \tilde{w}_i \right) = S_{p_k}$$

$$\delta \tilde{p}_s + \alpha \delta t \Pi_{1i} \delta D_i^* = S_{\tilde{p}_s}$$

To solve it, we put p_s and p into w to get w as function of D
 Then put w as function of D into h and p , thus we get p , h ,
 and p_s are all function of D , then put them into D equation
 to solve D

Hydrostatic IC

- Surface pressure
- Divergence/vorticity $\Rightarrow U/V$
- Use coordinate pressure to interpolation
- Use hydrostatic pressure as coordinate pressure and nonhydrostatic pressure
- Generate vertical velocity from $w=dz/dt$ equation

From hydrostatic system, we have all p at interfaces as

$$\hat{p}_k = \hat{A}_k + \hat{B}_k p_s \quad \& \quad p_k = \frac{1}{2}(\hat{p}_k + \hat{p}_{k+1})$$

From hydrostatic relation, we have

$$\hat{z}_{k+1} = \hat{z}_k + (\hat{p}_k - \hat{p}_{k+1}) \frac{\kappa_k h_k}{p_k g}$$

So coordinate pressure in hydrostatic system is

$$\left(\hat{\tilde{p}}_k \right)_{hydro} = \left(\hat{\tilde{p}}_{k+1} \right)_{hydro} + \frac{(a + \hat{z}_{k+1})^3 - (a + \hat{z}_k)^3}{3a^2} \left(\frac{p\bar{g}}{\kappa h} \right)_k$$

The new nonhydrostatic coordinate pressure is

$$\hat{\tilde{p}}_k = \hat{A}_k + \hat{B}_k \hat{\tilde{p}}_s$$

So interpolate u , v , h , and q from hydro to new coordinates

Remaining p and w at deep atmospheric system

From coordinate pressure definition

$$\hat{r}_{k+1}^3 = \hat{r}_k^3 + 3 \left(\frac{\kappa h}{p \bar{g}} \right)_k a^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right)$$

and hydrostatic relation

$$\hat{r}_{k+1} = \hat{r}_k + \left(\hat{p}_k - \hat{p}_{k+1} \right) \frac{\kappa_k h_k}{p_k \bar{g}}$$

combine both by eliminating $kh/(pg)$, we have

$$\hat{p}_k - \hat{p}_{k+1} = 3 \left(\frac{\hat{r}_{k+1} - \hat{r}_k}{\hat{r}_{k+1}^3 - \hat{r}_k^3} \right) a^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right)$$

can obtain p from top to surface

$$p_k = \frac{1}{2} \left(\hat{p}_k + \hat{p}_{k+1} \right)$$

Get vertical velocity w in deep atmospheric system

$$\text{from } \hat{r}_{k+1}^3 = \hat{r}_k^3 + 3 \left(\frac{\kappa h}{p \bar{g}} \right)_k a^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right) \text{ Do d/dt}$$

$$\text{We have } \hat{\tilde{w}}_{k+1} = \hat{\tilde{w}}_k + \frac{\kappa_k h_k}{p_k g} \Delta B_k \frac{d \hat{\bar{p}}_s}{dt} - \Delta \hat{\bar{p}}_k \frac{\kappa_k h_k}{\gamma p_k^2 g} \frac{dp_k}{dt}$$

where

$$\frac{d \hat{\bar{p}}_s}{dt} = - \sum_{k=1}^K m^2 \left(\hat{\bar{p}}_k - \hat{\bar{p}}_{k+1} \right) \left(\frac{\partial u^*}{\partial \lambda} + \frac{\partial v^*}{\partial \varphi} \right)_k$$

$$\left(\frac{dp}{dt} \right)_k = -\gamma_k p_k \left(\nabla_H \cdot (V_H) + \frac{1}{\Delta r^3} V_H \cdot \nabla_H \Delta r^3 - \frac{1}{\Delta r^3} \Delta (V_H \cdot \nabla_H r^3) + \frac{a^2}{r^2} \frac{\Delta \tilde{w}}{\Delta r} \right)_k$$

$$\text{and } \hat{\tilde{w}}_1 = m^2 \frac{\hat{r}_1}{a} \left(u_1^* \frac{\partial \hat{r}_1}{a \partial \lambda} + v_1^* \frac{\partial \hat{r}_1}{a \partial \varphi} \right) \quad \& \quad \tilde{w}_k = \frac{1}{2} \hat{\tilde{w}}_{k+1} + \frac{1}{2} \hat{\tilde{w}}_k$$

Summary

- Deep atmospheric dynamics is discretized in generalized hybrid coordinates
- Semi-implicit semi-Lagrangian is used
- Initial condition with hydrostatic system is used
- It is suitable for very high resolution and coupling with space model
- Details can be found in Juang 2014, NCEP office note #477
- Coded and still under testing.