

Free and convectively coupled equatorial waves diagnosis using 3-D Normal Modes

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Convectively coupled equatorial-wave diagnosis using three-dimensional normal modes

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A new methodology for the diagnosis of convectively coupled equatorial waves (CCEWs) is presented. It is based on a pre-filtering of the geopotential and horizontal wind, using three-dimensional (3D) normal mode functions of the adiabatic linearized equations of a resting atmosphere, followed by a space–time spectral analysis to identify the spectral regions of coherence.

The methodology permits a direct detection of various types of equatorial wave, compares the dispersion characteristics of the coupled waves with the theoretical dispersion curves and allows an identification of which vertical modes are more involved in the convection. Moreover, the proposed methodology is able to show the existence of free dry waves and moist coupled waves with a common vertical structure, which is in conformity with the effect of convective heating/cooling on the effective static stability, as deduced from the

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Outline

Observation of Convectively Coupled Equatorial Waves (CCEW)

Theory of equatorial waves

Data and Method

Results

- Coherence spectra

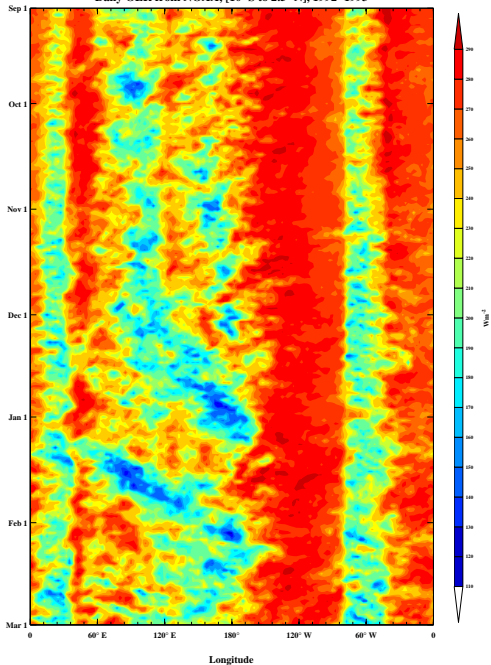
- Vertical decomposition

- Power spectra

- Separating stratosphere and troposphere

- QBO effects on the vertical propagation of equatorial waves

Conclusions

Daily OLR from NOAA, [10° S to 2.5° N], 1992–1993

Equatorial waves

- ▶ Linearized primitive equations:

$$\frac{\partial u}{\partial t} - f v + \frac{1}{a \cos \theta} \frac{\partial \phi}{\partial \lambda} = 0 \quad (1)$$

$$\frac{\partial v}{\partial t} + f u + \frac{1}{a} \frac{\partial \phi}{\partial \theta} = 0 \quad (2)$$

$$\nabla \cdot \mathbf{V} + \frac{1}{\rho_0} \frac{\partial}{\partial z} (\rho_0 w) = 0 \quad (3)$$

$$\frac{1}{\rho_0} \frac{\partial}{\partial t} \left(\frac{\partial \phi}{\partial z} \right) + w N^2 = 0 \quad (4)$$

Equatorial waves

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- ▶ Assuming N constant and combining (3) and (4) to eliminate w :

Equatorial waves

- ▶ Linearized primitive equations:

$$\frac{\partial u}{\partial t} - f v + \frac{1}{a \cos \theta} \frac{\partial \phi}{\partial \lambda} = 0 \quad (5)$$

$$\frac{\partial v}{\partial t} + f u + \frac{1}{a} \frac{\partial \phi}{\partial \theta} = 0 \quad (6)$$

$$\frac{\partial}{\partial t} \left[\frac{1}{\rho_0} \frac{\partial}{\partial z} \left(\rho_0 \frac{\partial \phi}{\partial z} \right) \right] - N^2 \nabla \cdot \mathbf{v} = 0 \quad (7)$$

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- ▶ Considering solutions with separable vertical structure:

$$[u, v, \phi] = G(z) \left[\tilde{u}(t, \theta, \lambda), \tilde{v}(t, \theta, \lambda), \tilde{\phi}(t, \theta, \lambda) \right]$$

Equatorial waves

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- ▶ Like in a wave

$$[u, v, \phi] = \left[\hat{u}(\theta), \hat{v}(\theta), \hat{\phi}(\theta) \right] e^{z/2H} e^{imz} e^{i(k\lambda + \omega t)}$$

Equatorial waves

- ▶ One obtains the horizontal structure equations:

$$\frac{\partial \tilde{u}}{\partial t} - f \tilde{v} + \frac{1}{a \cos \theta} \frac{\partial \tilde{\phi}}{\partial \lambda} = 0 \quad (8)$$

$$\frac{\partial \tilde{v}}{\partial t} + f \tilde{u} + \frac{1}{a} \frac{\partial \tilde{\phi}}{\partial \theta} = 0 \quad (9)$$

$$\frac{\partial \tilde{\phi}}{\partial t} + gh_e \nabla \cdot \tilde{\mathbf{V}} = 0 \quad (10)$$

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- ▶ and the vertical structure equation:

$$\frac{1}{\rho_0} \frac{\partial}{\partial z} \left(\rho_0 \frac{\partial G}{\partial z} \right) + \frac{N^2}{gh_e} G = 0 \quad (11)$$

- ▶ where $-1/gh_e$ is the separation constant.

Equatorial waves

- ▶ For a wave with vertical wave number m :

$$[u, v, \phi] = \left[\hat{u}(\theta), \hat{v}(\theta), \hat{\phi}(\theta) \right] e^{z/2H} e^{imz} e^{i(k\lambda + \omega t)}$$

Equatorial waves

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$$[u, v, \phi] = \left[\hat{u}(\theta), \hat{v}(\theta), \hat{\phi}(\theta) \right] e^{z/2H} e^{imz} e^{i(k\lambda + \omega t)}$$

- ▶ the equivalent depth is given by

$$h_e = \frac{N^2}{g \left(m^2 + \frac{1}{4H^2} \right)} \quad (12)$$

Equatorial waves

- ▶ Using the equatorial β -plane approximation:

$$\frac{\partial \tilde{u}}{\partial t} - \beta y \tilde{v} + \frac{\partial \tilde{\phi}}{\partial x} = 0 \quad (13)$$

$$\frac{\partial \tilde{v}}{\partial t} + \beta y \tilde{u} + \frac{\partial \tilde{\phi}}{\partial y} = 0 \quad (14)$$

$$\frac{\partial \tilde{\phi}}{\partial t} + gh_e \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0 \quad (15)$$

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- ▶ A complete set of zonal propagating wave solutions of these system of shallow water equations was found by *Matsuno* (1966).

3-D normal mode basis

- ▶ We solved the equations over the sphere in isobaric coordinates, with the vertical structure equation given by:

$$\frac{\partial}{\partial p} \left(\frac{1}{S_0} \frac{\partial G}{\partial p} \right) + \frac{1}{gh_e} G = 0$$

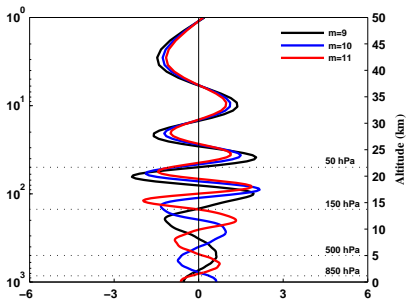
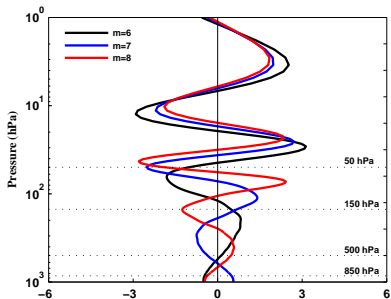
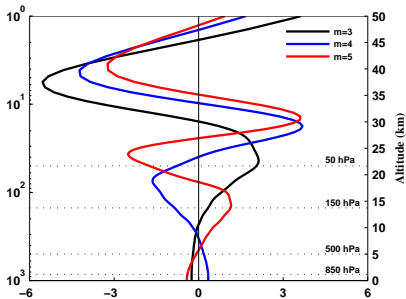
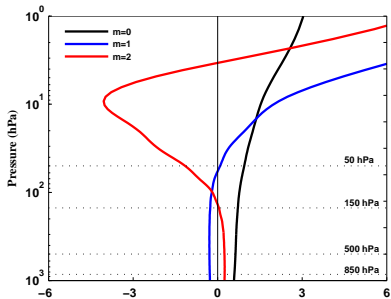
3-D normal mode basis

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$$\frac{\partial}{\partial p} \left(\frac{1}{S_0} \frac{\partial G}{\partial p} \right) + \frac{1}{gh_e} G = 0$$

- ▶ The VSE was solved numerically with a spectral method as in *Kasahara (1984)* and *Castanheira et al. (1999)*.

Vertical Structure Functions



3-D normal mode basis

- ▶ The horizontal structure equations over the sphere

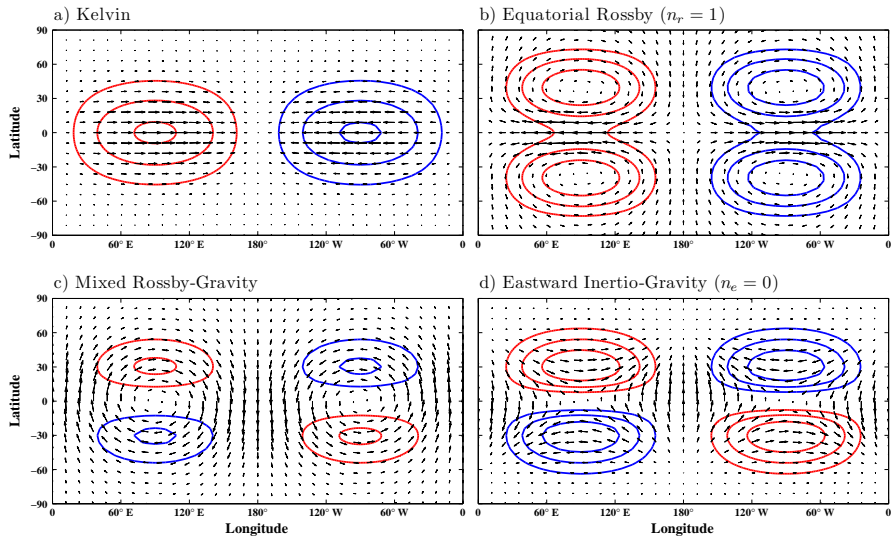
$$\frac{\partial \tilde{u}}{\partial t} - f \tilde{v} + \frac{1}{a \cos \theta} \frac{\partial \tilde{\phi}}{\partial \lambda} = 0 \quad (16)$$

$$\frac{\partial \tilde{v}}{\partial t} + f \tilde{u} + \frac{1}{a} \frac{\partial \tilde{\phi}}{\partial \theta} = 0 \quad (17)$$

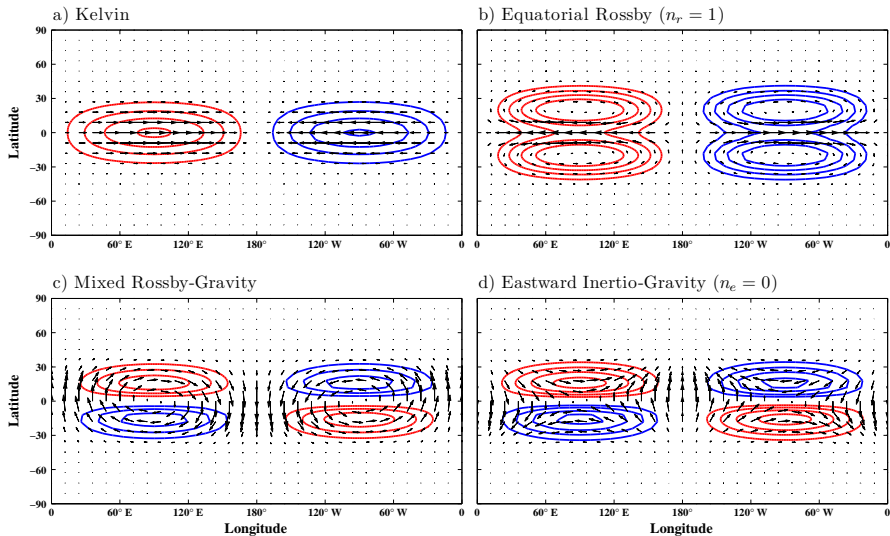
$$\frac{\partial \tilde{\phi}}{\partial t} + gh_e \nabla \cdot \tilde{\mathbf{V}} = 0 \quad (18)$$

- ▶ were solved using the methodology of *Swarztrauber and Kasahara (1985)*.

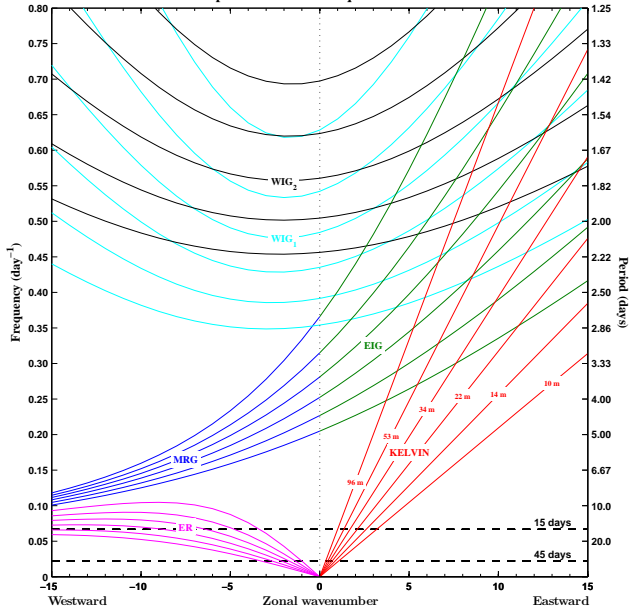
Shallow water waves over the sphere ($h_1 = 5591$ m)



Shallow water waves over the sphere ($h_4 = 474$ m)



Dispersion curves for equatorial waves



Data

- ▶ Outgoing Longwave Radiation (OLR) from the National Oceanic and Atmospheric Administration (NOAA) for the period 1979-2012.
- ▶ Horizontal wind (u, v) and geopotential (ϕ) from the ERA-Interim reanalysis (1979-2012).

Method: filtering the dynamical fields

- Projection of the horizontal wind (u, v) and geopotential (ϕ) onto the normal modes of the linearized primitive equations on the sphere.

$$\begin{bmatrix} u \\ v \\ \phi \end{bmatrix} = \sum_{m,k,n} w_{mkn}(t) G_m(p) e^{ik\lambda} \mathbf{C}_m \cdot \begin{bmatrix} U(\theta) \\ iV(\theta) \\ Z(\theta) \end{bmatrix}_{mkn}$$

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- ▶ Reconstitution of the horizontal wind (u, v) and geopotential (ϕ) with a given subset of modes (filtering):

$$\begin{bmatrix} u_n \\ v_n \\ \phi_n \end{bmatrix} = \sum_{m,k} w_{mkn}(t) G_m(p) e^{ik\lambda} \mathbf{C}_m \cdot \begin{bmatrix} U(\theta) \\ iV(\theta) \\ Z(\theta) \end{bmatrix}_{mkn}$$

Method: space-time cross-spectral analysis (Hayashi, 1982)

- ▶ Considering a space-time series

$$Y(\lambda, t) = \sum_k [C_k(t) \cos(k\lambda) + S_k(t) \sin(k\lambda)]$$

Method: space-time cross-spectral analysis (Hayashi, 1982)

- ▶ Considering a space-time series

$$Y(\lambda, t) = \sum_k [C_k(t) \cos(k\lambda) + S_k(t) \sin(k\lambda)]$$

- ▶ The space-time power spectra is given by

$$4 P_{k, \pm \omega}(Y) = P_{\omega}(C_k) + P_{\omega}(S_k) \pm 2 Q_{\omega}(C_k, S_k),$$

where P_{ω} and Q_{ω} are the time power and quadrature spectra, respectively.

Method: space-time cross-spectral analysis (Hayashi, 1982)

- ▶ Spectral coherence and phase difference between two fields $Y(\lambda, t)$ and $Y'(\lambda, t)$ are given by

$$\text{Coh}_{k,\pm\omega}^2(Y, Y') = \frac{K_{k,\pm\omega}^2(Y, Y') + Q_{k,\pm\omega}^2(Y, Y')}{P_{k,\pm\omega}(Y)P_{k,\pm\omega}(Y')}$$

- ▶ and

$$\text{Ph}_{k,\pm\omega}(Y, Y') = \tan^{-1} [Q_{k,\pm\omega}(Y, Y')/K_{k,\pm\omega}(Y, Y')]$$

- ▶ where

Method: space-time cross-spectral analysis (Hayashi, 1982)

- ▶ where

$$4 K_{k,\pm\omega}(Y, Y') = K_{\omega}(C_k, C_k') + K_{\omega}(S_k, S_k') \\ \pm Q_{\omega}(C_k, S_k') \mp Q_{\omega}(S_k, C_k')$$

- ▶ and

$$4 Q_{k,\pm\omega}(Y, Y') = \pm Q_{\omega}(C_k, C_k') \pm Q_{\omega}(S_k, S_k') \\ - K_{\omega}(C_k, S_k') + K_{\omega}(S_k, C_k')$$

- ▶ are cospectra and quadrature spectra, respectively.

Testing the methodology

- ▶ Reconstructing the circulation field

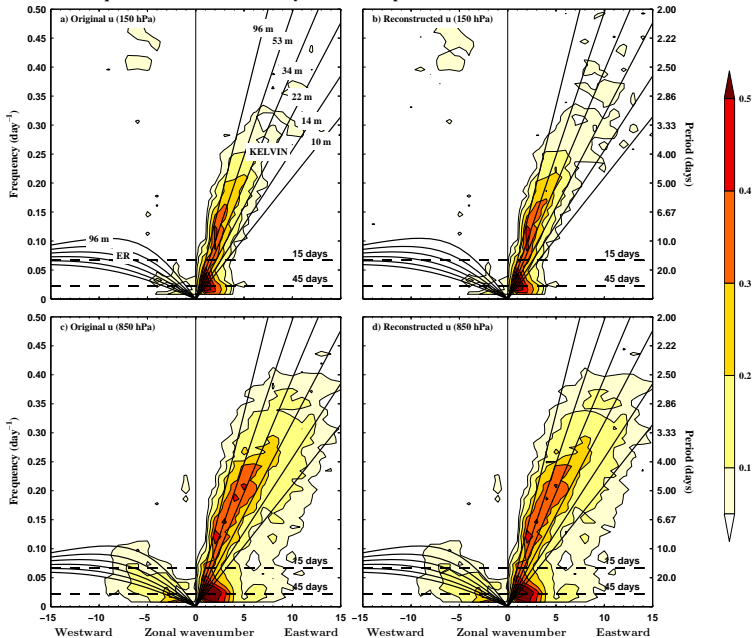
$$\begin{bmatrix} u \\ v \\ \phi \end{bmatrix} = \sum_{m,k,n} w_{mkn}(t) G_m(\rho) e^{ik\lambda} \mathbf{C}_m \cdot \begin{bmatrix} U(\theta) \\ iV(\theta) \\ Z(\theta) \end{bmatrix}_{mkn}$$

- ▶ Symmetric and antisymmetric components of a variable Y

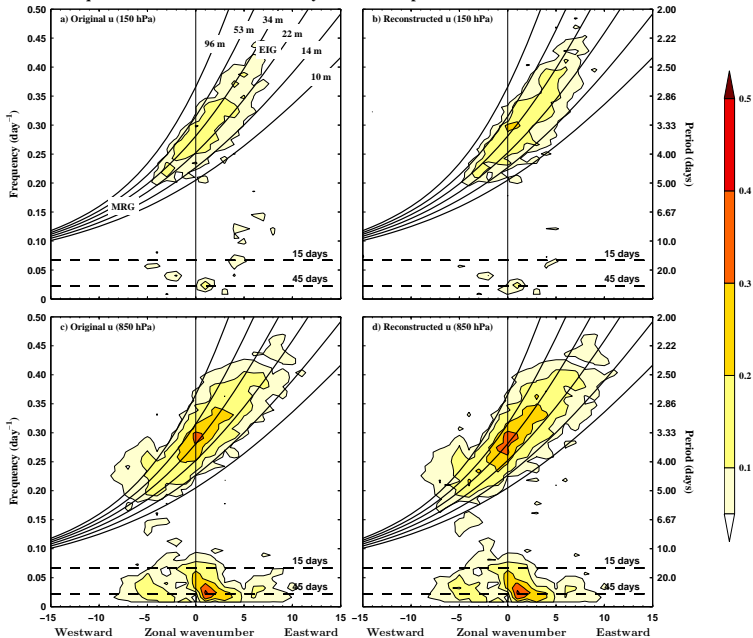
$$Y_S(\theta) = \frac{Y(\theta) + Y(-\theta)}{2}$$

$$Y_A(\theta) = \frac{Y(\theta) - Y(-\theta)}{2}$$

Squared coherence between symmetric components of zonal wind and OLR



Squared coherence between anti-symmetric components of zonal wind and OLR



Selecting waves

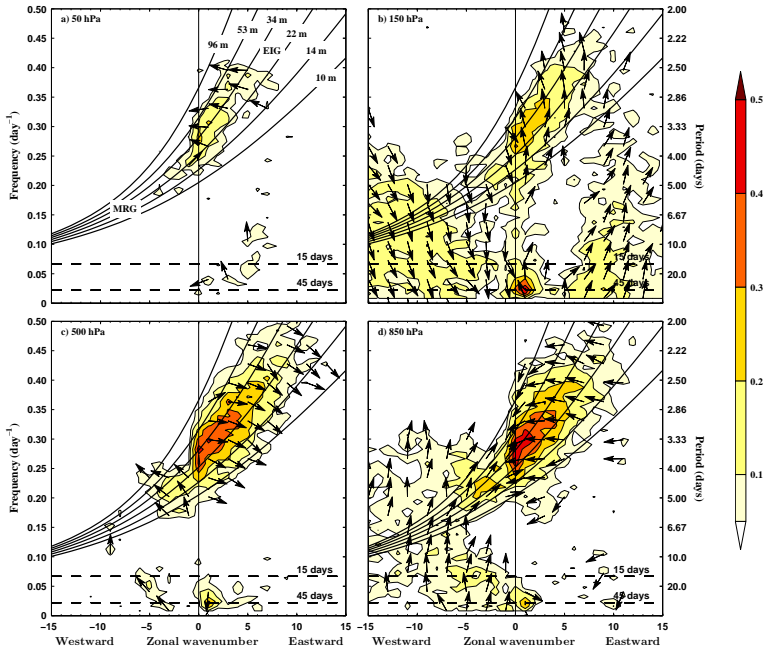
- ▶ Fixing a given pair of zonal, k and meridional n indices, waves of a given type are selected

$$\begin{bmatrix} u_{kn} \\ v_{kn} \\ \phi_{kn} \end{bmatrix} = \sum_m w_{mkn}(t) G_m(\rho) e^{ik\lambda} \mathbf{C}_m \cdot \begin{bmatrix} U(\theta) \\ iV(\theta) \\ Z(\theta) \end{bmatrix}_{mkn}$$

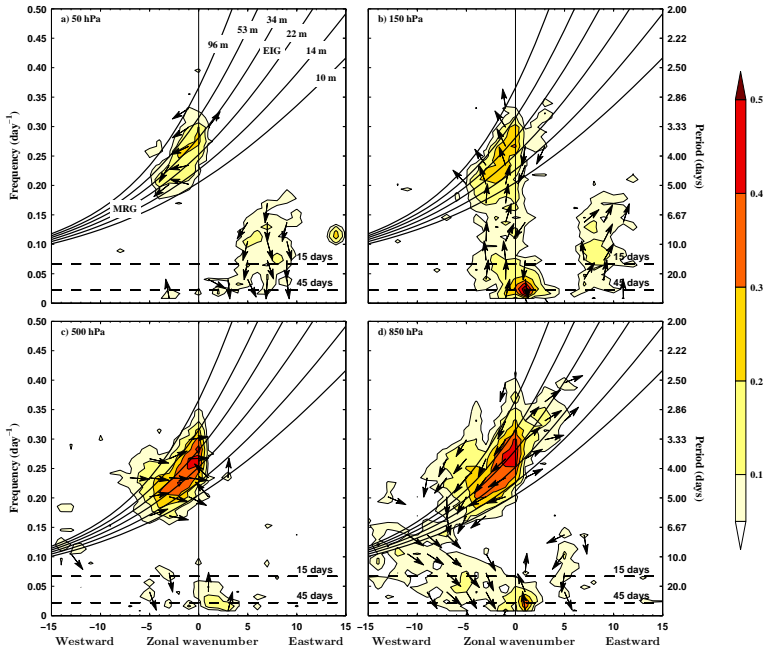
Results

- ▶ Space-time coherence spectra for eastward inertio-gravity (EIG), mixed Rossby-gravity (MRG), Kelvin (Kel) and Equatorial Rossby (ER) waves.

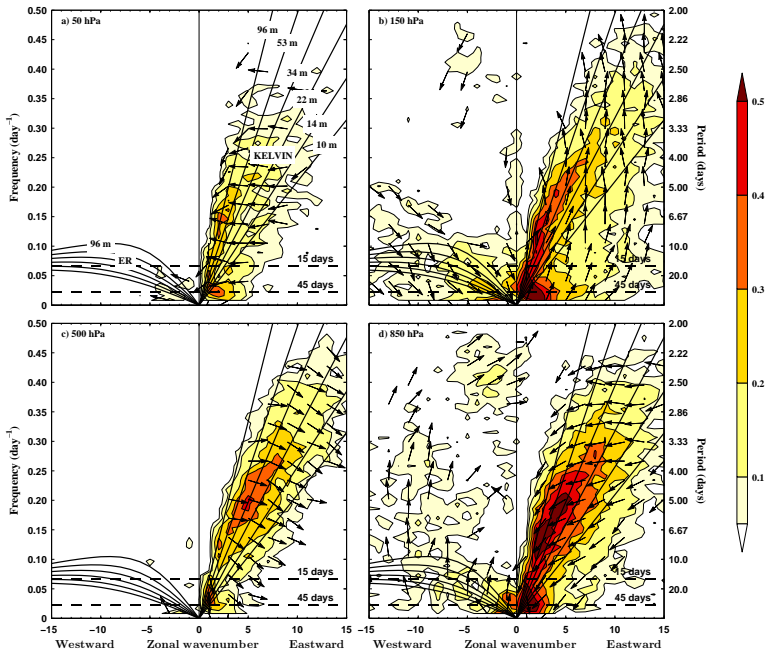
Squared coherence between u_{EIG} and anti-symmetric OLR



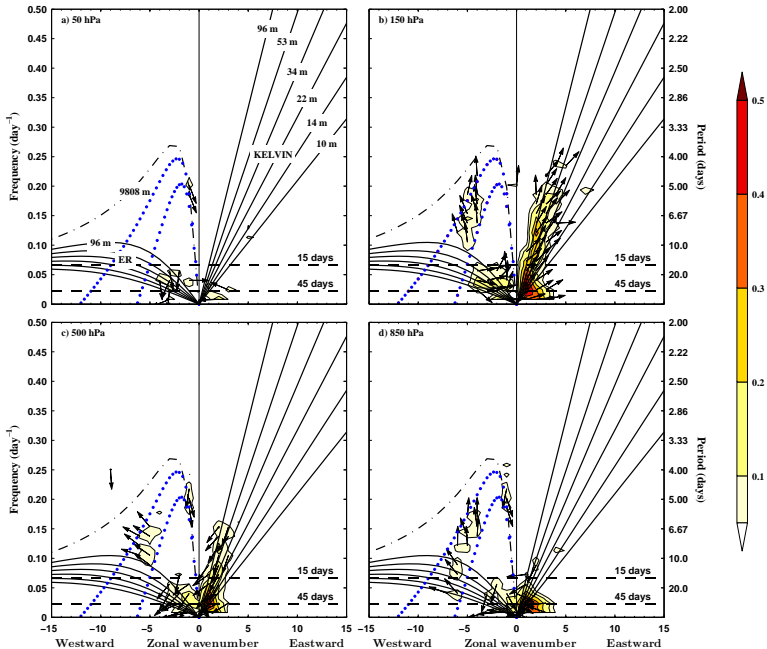
Squared coherence between u_{MRG} and anti-symmetric OLR



Squared coherence between u_{KEL} and symmetric OLR



Squared coherence between u_{ER} and symmetric OLR



Vertical mode decomposition of waves

- ▶ Fixing the vertical, zonal, and meridional indices, mkn , respectively, we obtain a vertical modal decomposition of the waves

$$\begin{bmatrix} u_{mkn} \\ v_{mkn} \\ \phi_{mkn} \end{bmatrix} = w_{mkn}(t) G_m(p) e^{ik\lambda} \mathbf{C} \cdot \begin{bmatrix} U(\theta) \\ iV(\theta) \\ Z(\theta) \end{bmatrix}_{mkn}$$

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- ▶ Now for, the cosine and sine coefficients are given by

$$C_k^{mn} \propto \Re [w_{mkn}(t)] \quad S_k^{mn} \propto \Im [w_{mkn}(t)]$$

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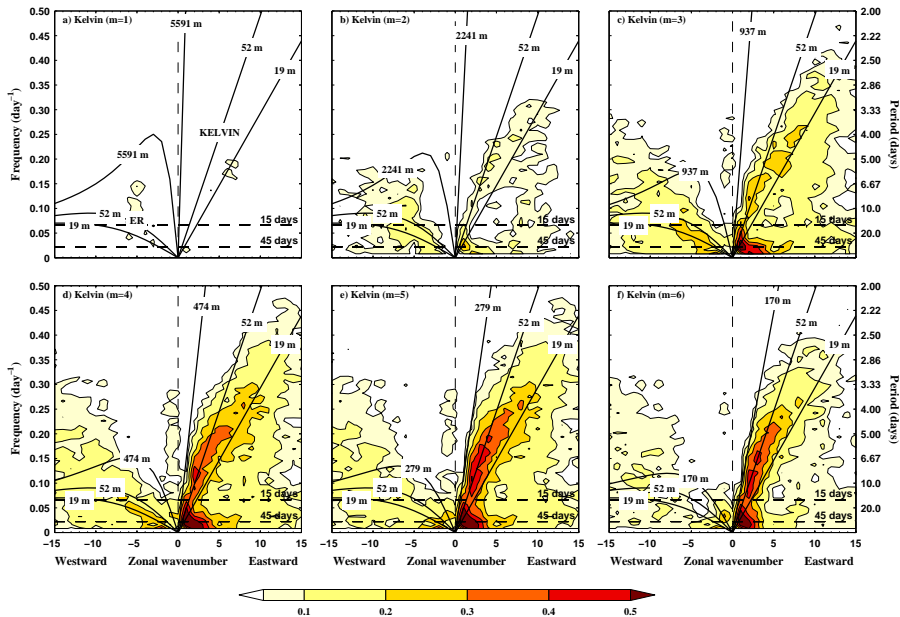
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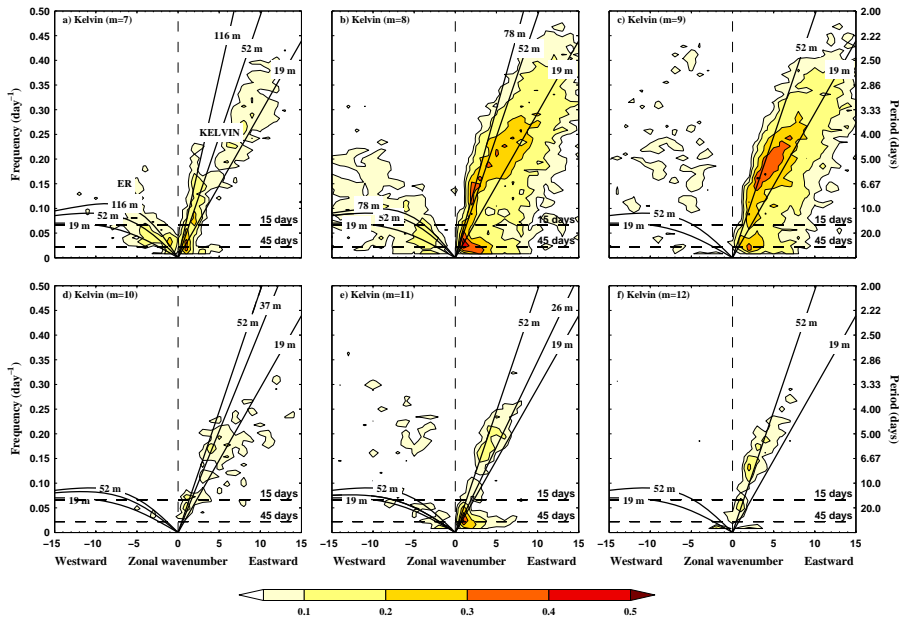
$$C_k^{mn} \propto \Re [w_{mkn}(t)] \quad S_k^{mn} \propto \Im [w_{mkn}(t)]$$

- ▶ and we may analyze the coherence of the 3-D waves with the OLR and interpret their power spectra as the total (Kinetic + Available Potential) energy spectra.

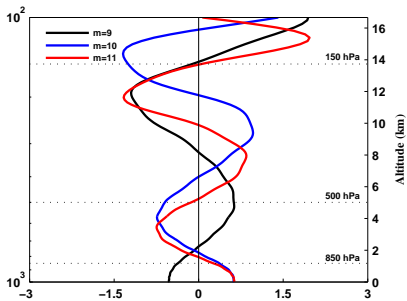
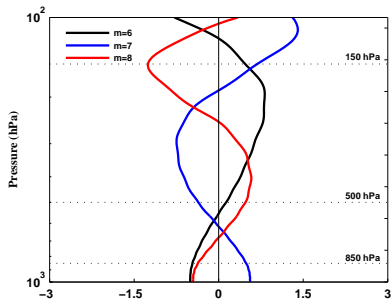
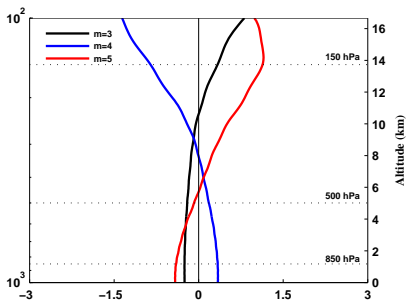
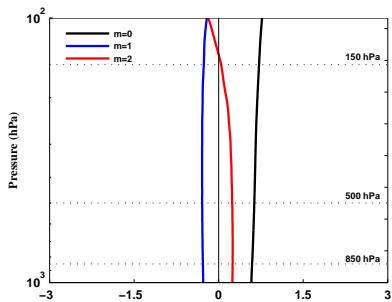
Squared coherence between the Kelvin baroclinic modes and symmetric OLR



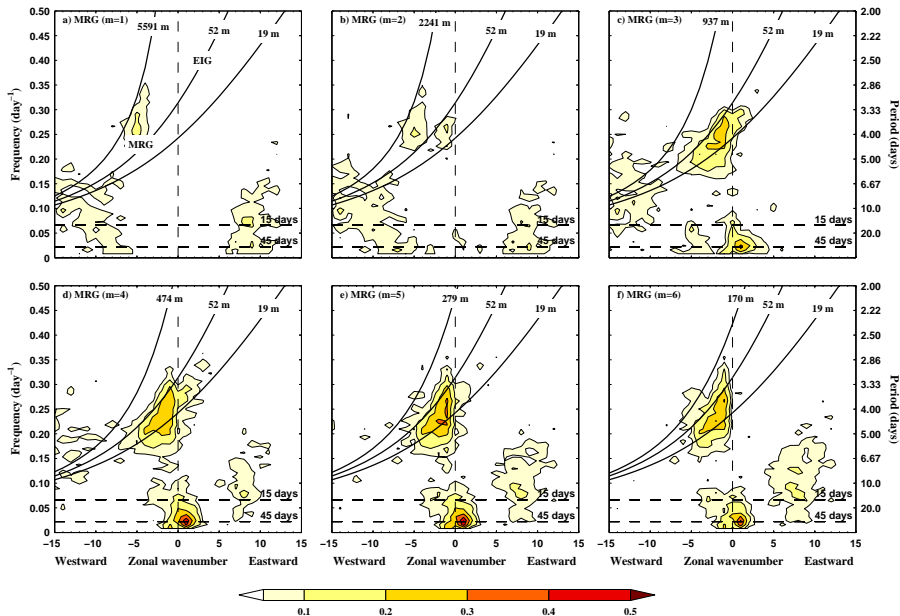
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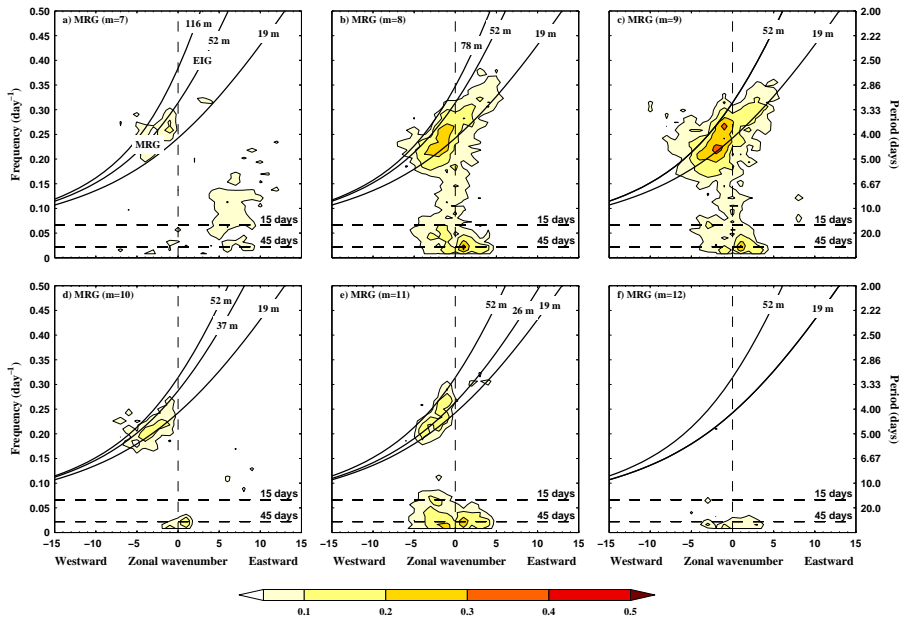
Vertical Structure Functions



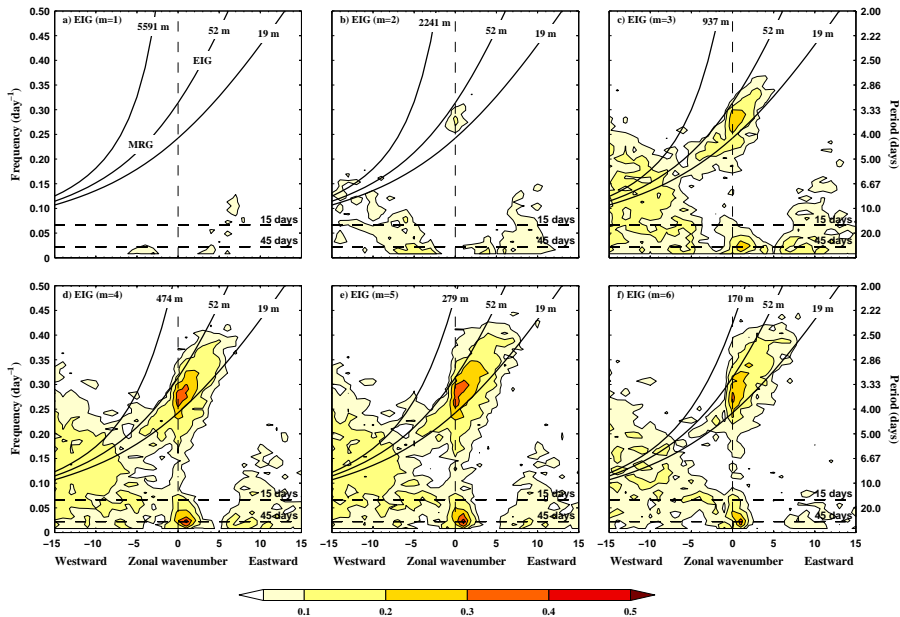
Squared coherence between the MRG baroclinic modes and anti-symmetric OLR



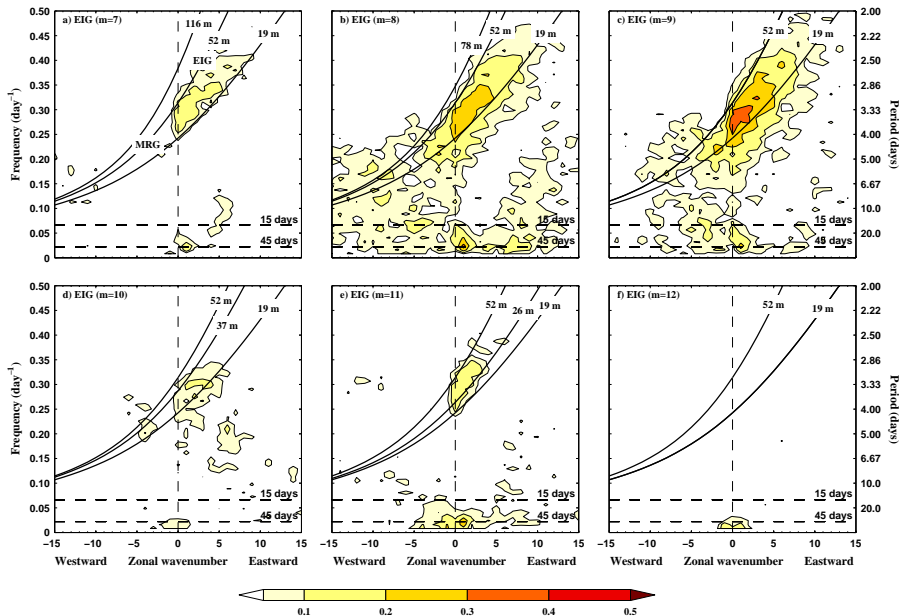
Squared coherence between the MRG baroclinic modes and anti-symmetric OLR



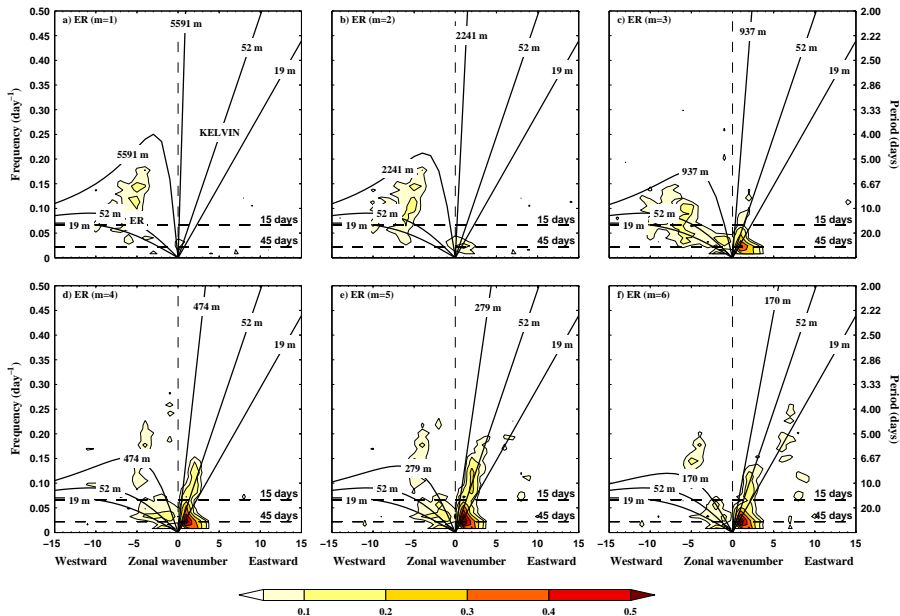
Squared coherence between the EIG baroclinic modes and anti-symmetric OLR



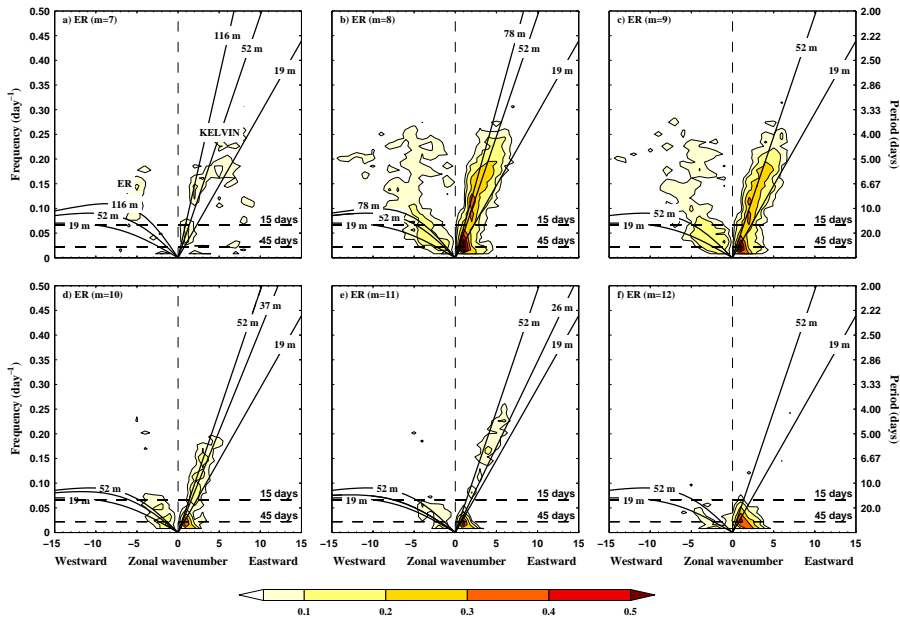
Squared coherence between the EIG baroclinic modes and anti-symmetric OLR



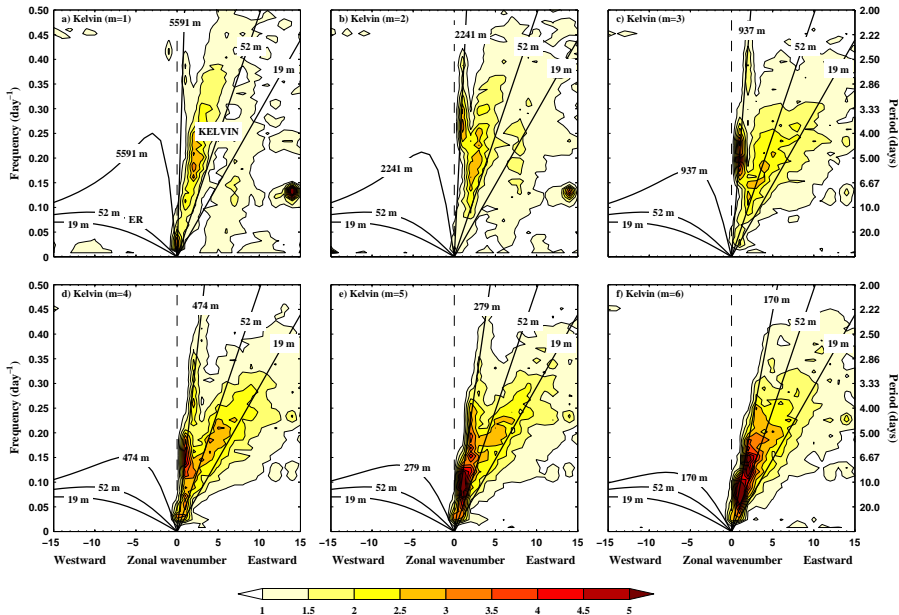
Squared coherence between the ER baroclinic modes and symmetric OLR



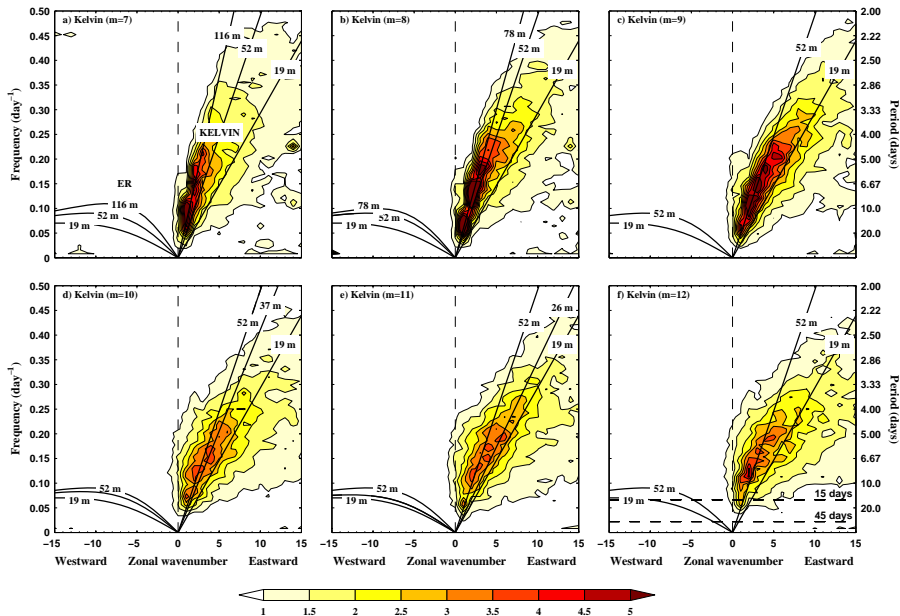
Squared coherence between the ER baroclinic modes and symmetric OLR



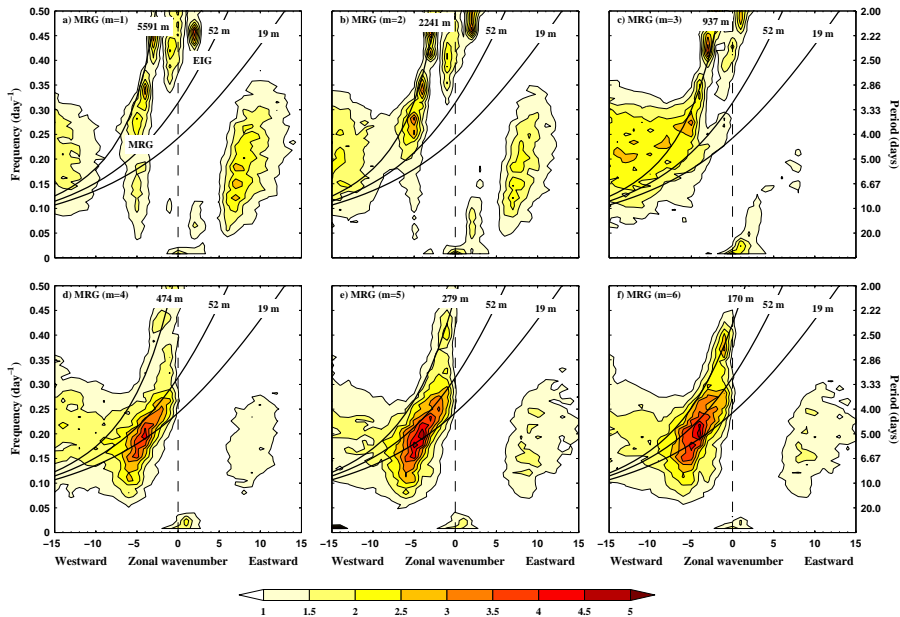
Normal Mode energy spectrum divided by the background red noise



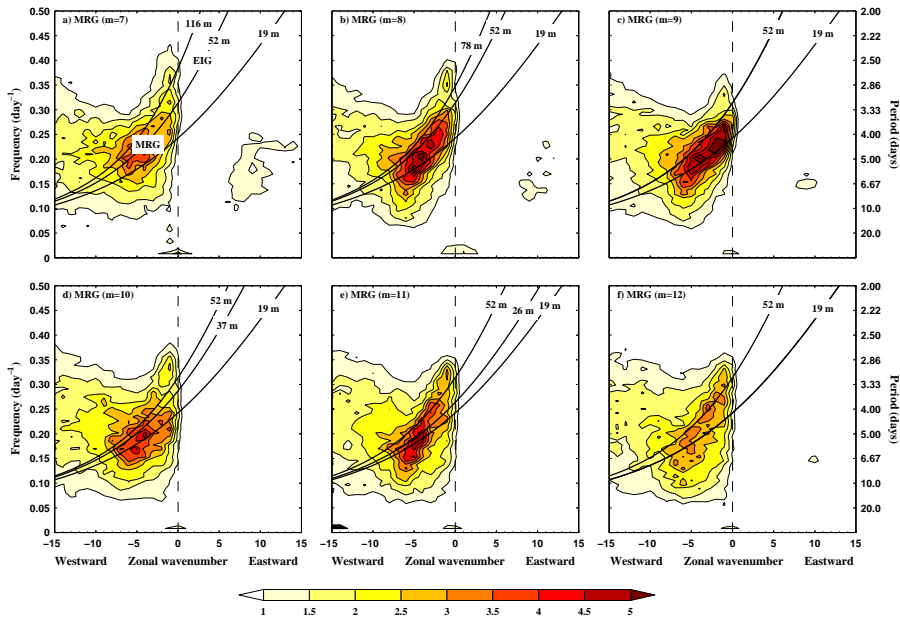
Normal Mode energy spectrum divided by the background red noise



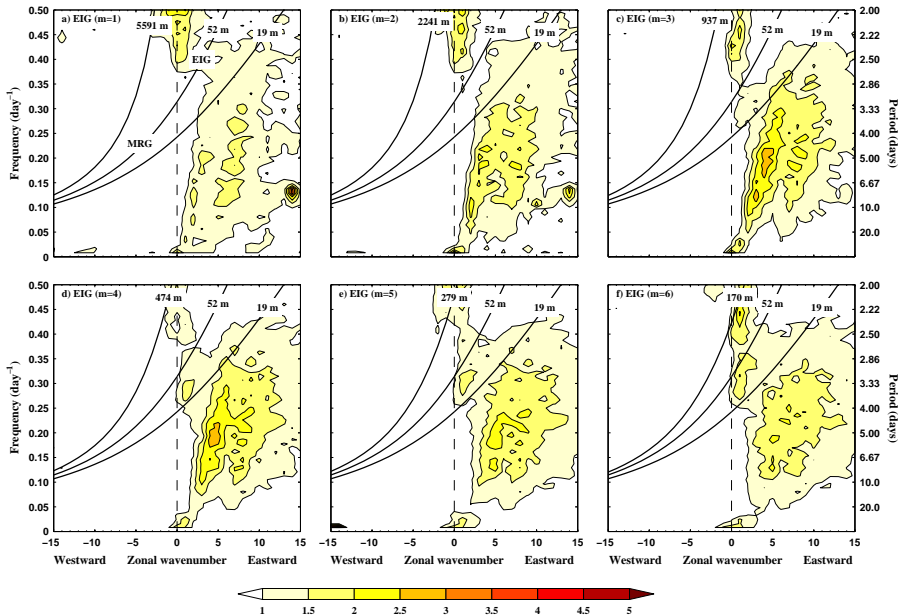
Normal Mode energy spectrum divided by the background red noise



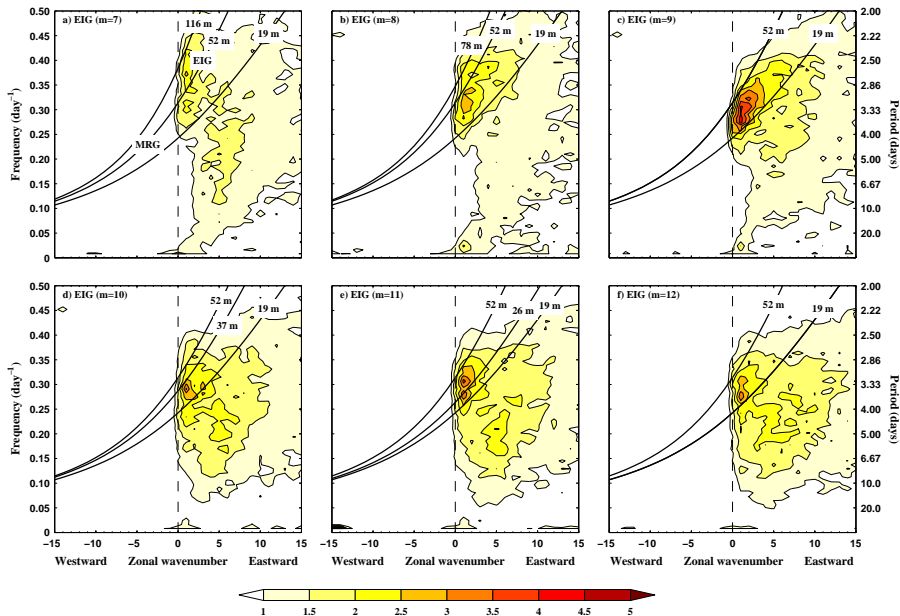
Normal Mode energy spectrum divided by the background red noise



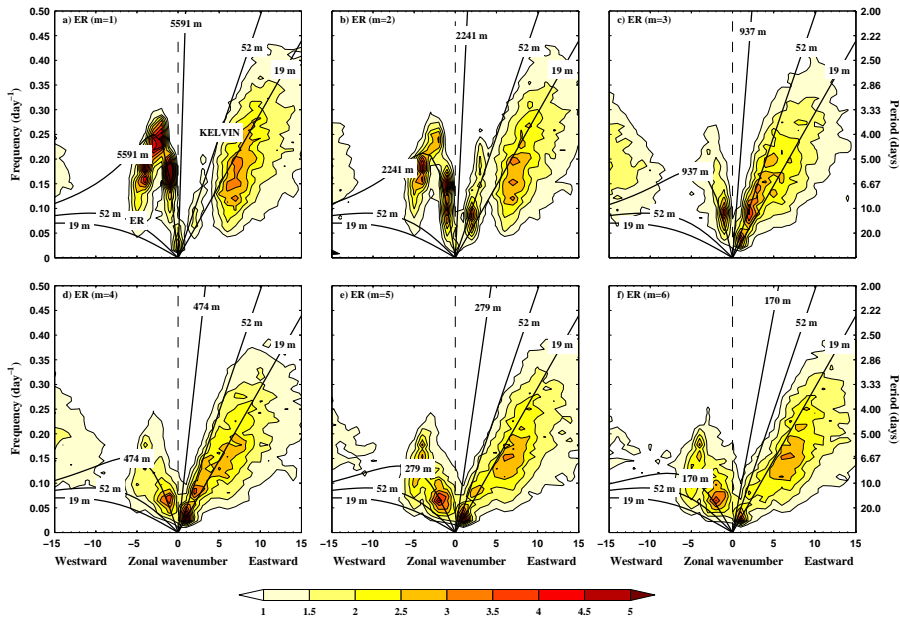
Normal Mode energy spectrum divided by the background red noise



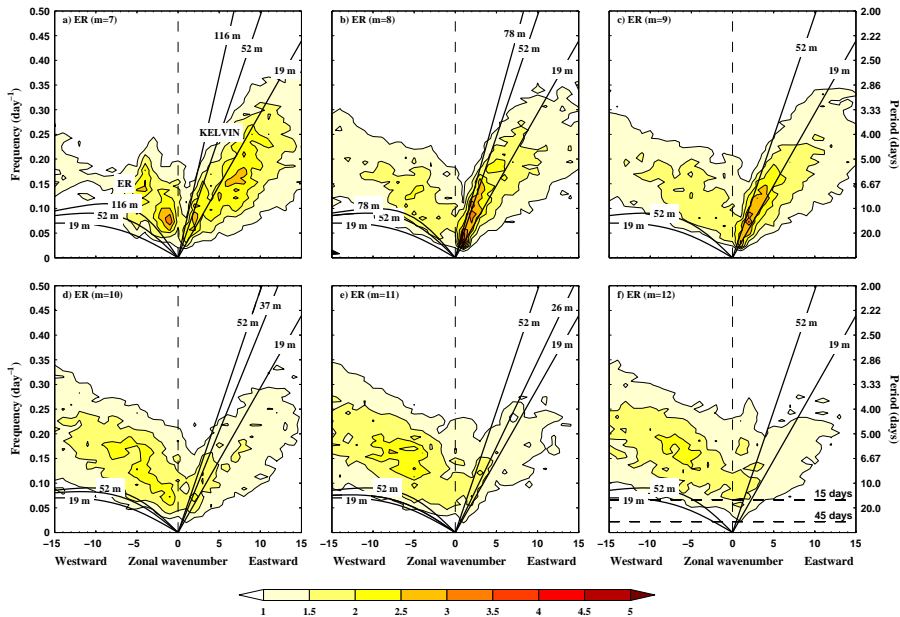
Normal Mode energy spectrum divided by the background red noise



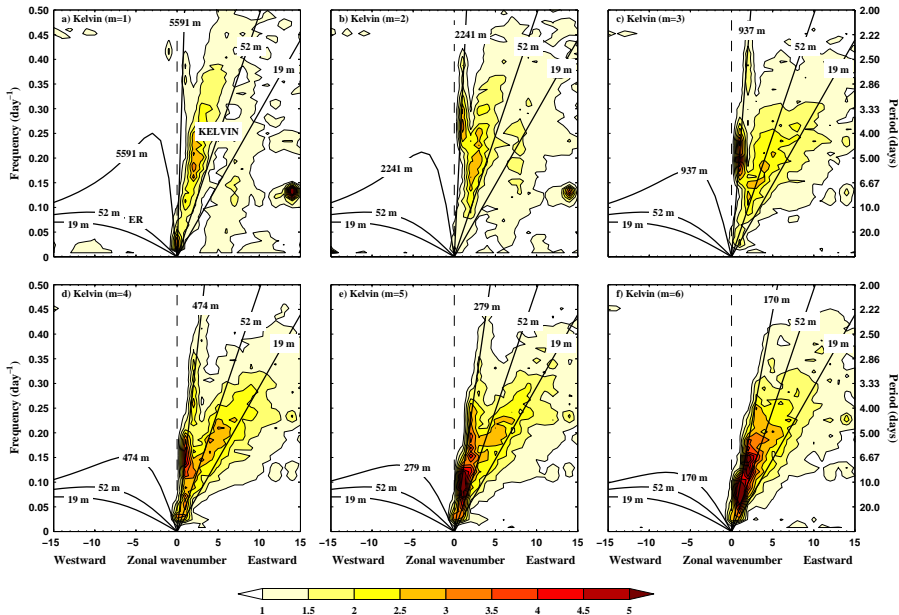
Normal Mode energy spectrum divided by the background red noise



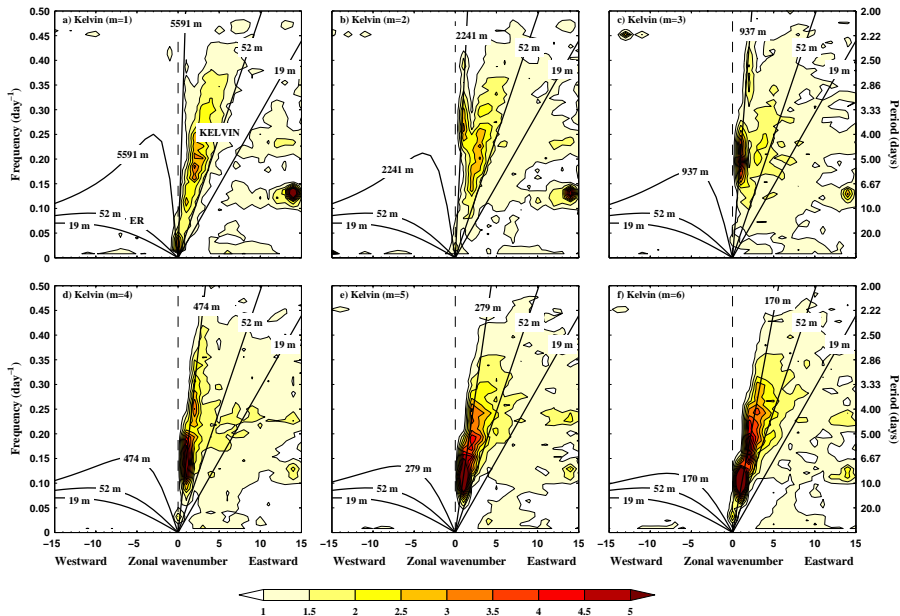
Normal Mode energy spectrum divided by the background red noise



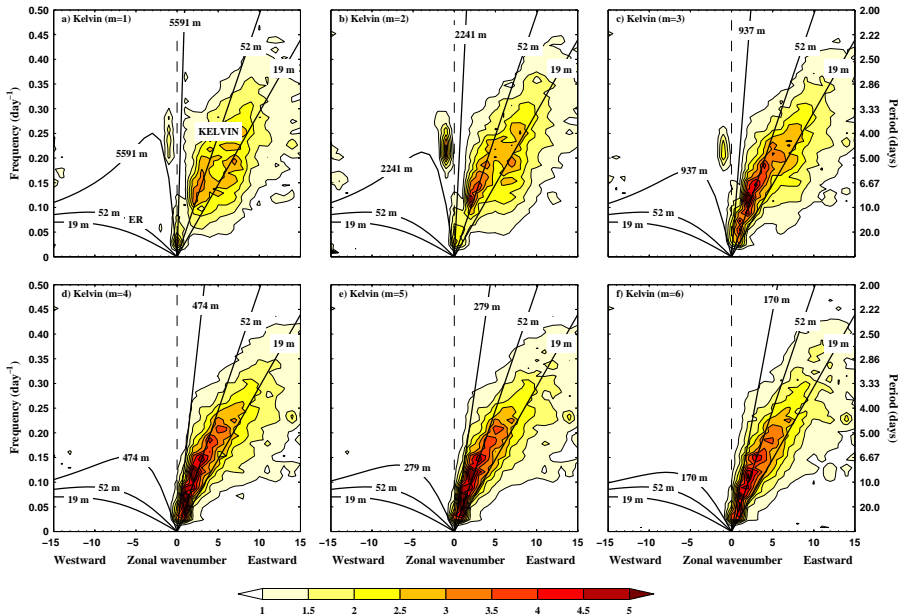
Normal Mode energy spectrum divided by the background red noise



Stratosphere : Normal Mode energy spectrum divided by the background red noise



Troposphere : Normal Mode energy spectrum divided by the background red noise



QBO effects

- ▶ The vertical propagation of equatorial waves was analysed for the two phases of the QBO. Considering the daily zonal mean 30-hPa zonal wind at the equator the QBO phases were defined as follows

$$\bar{u}(0 \text{ N}, 30 \text{ hPa}) > 5 \text{ m s}^{-1} \Rightarrow \textit{Westerly phase}$$

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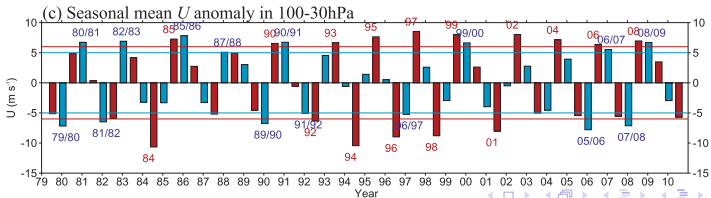
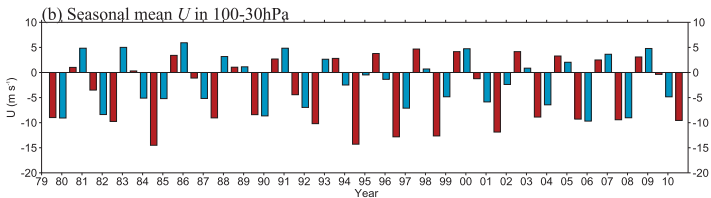
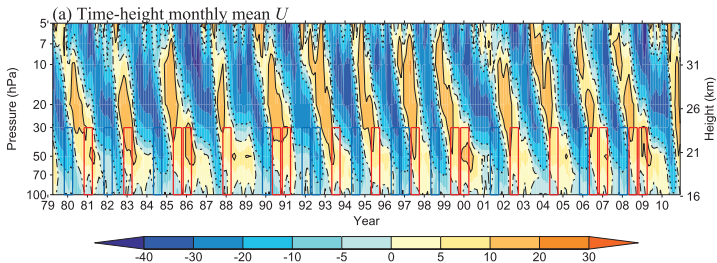
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- ▶ From the linear theory and using the slow-variation WKBJ approximation (*Andrews et al.*, 1987), waves with zonal phase velocity c can propagate vertically only if the zonal winds satisfy

$$c - \bar{u}(z) > 0 \quad \textit{for Kelvin waves}$$

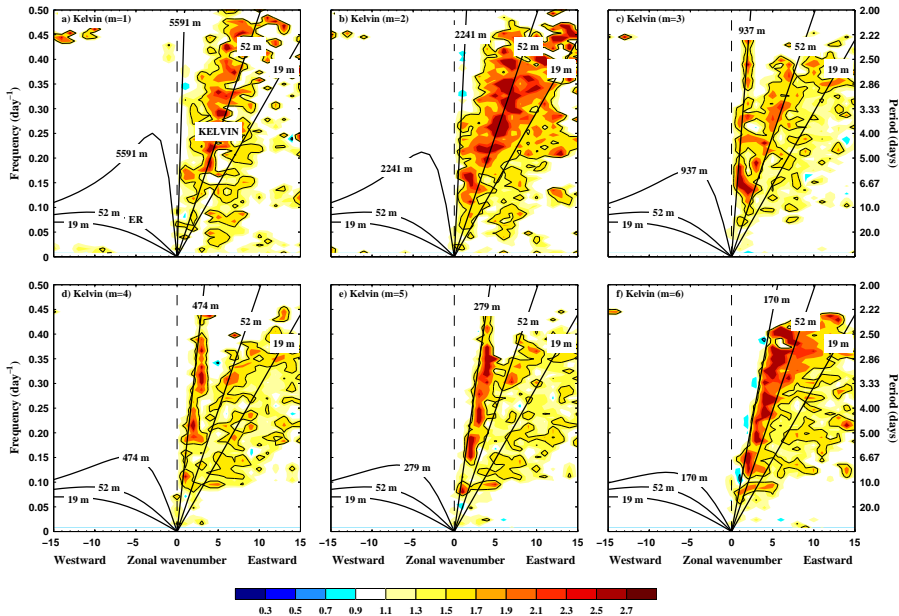
$$-\beta/k^2 < c - \bar{u}(z) < 0 \quad \textit{for MRG waves}$$



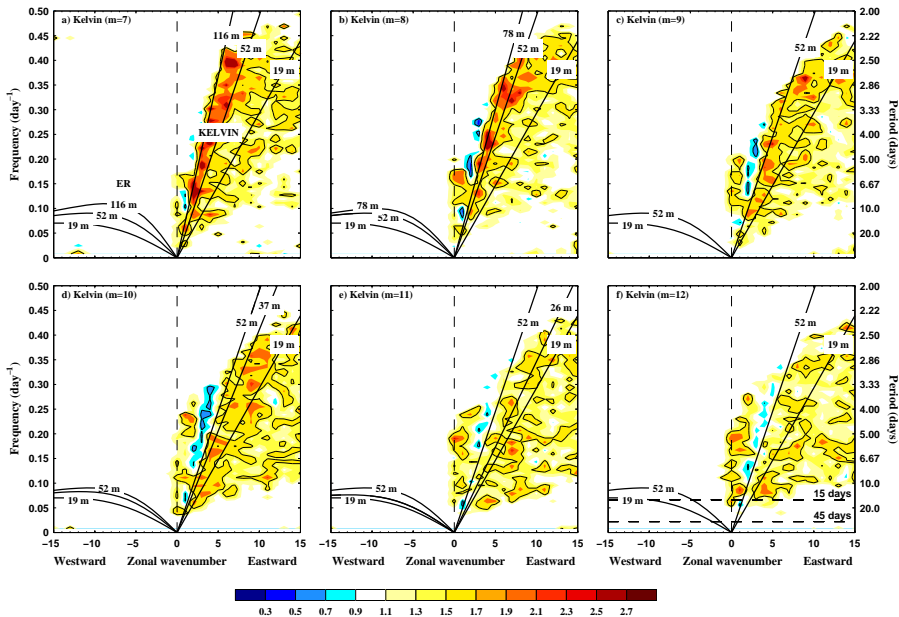
- ▶ The ratios between the power spectra calculated for each QBO phase are represented in the next Figures

$$F = \frac{P^{E-QBO}}{P^{W-QBO}}$$

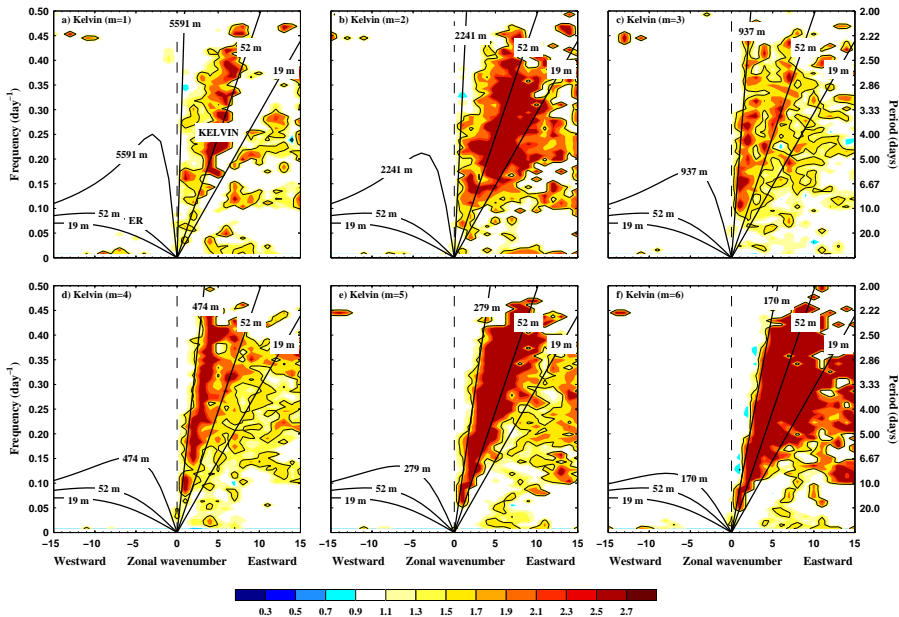
Ratio between the Normal Mode Power spectra for the two QBO phases (Easterly/Westerly)



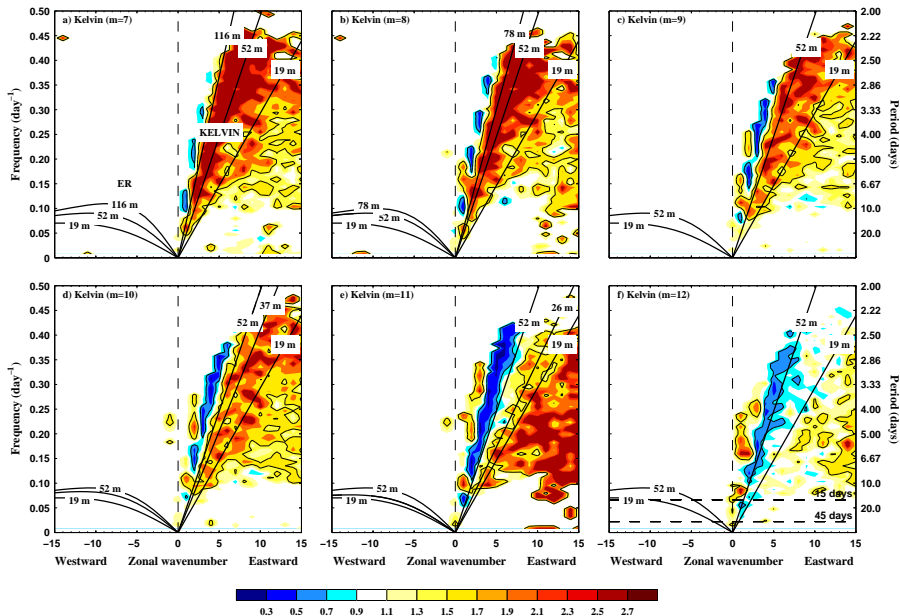
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Stratosphere : Ratio between the Normal Mode Power spectra for the two QBO phases (Easterly/Westerly)



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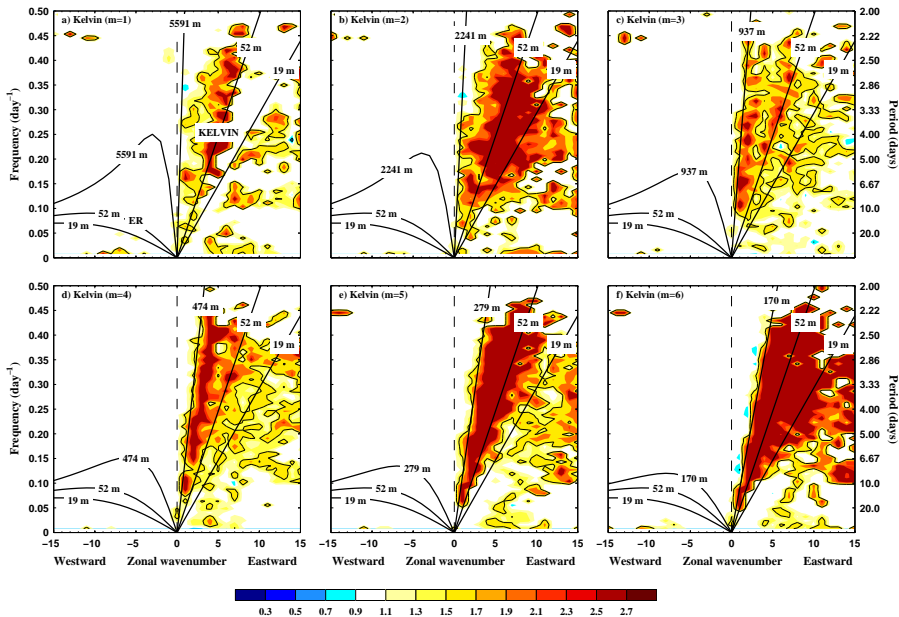
$$m(z) = \mp \frac{N}{c - \bar{u}}$$

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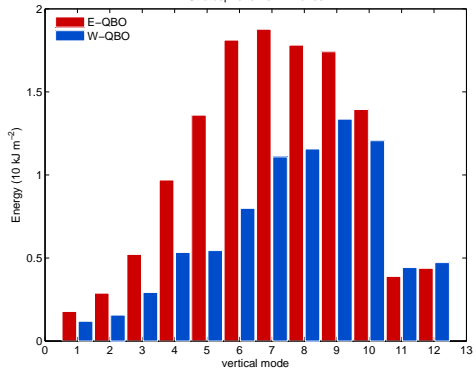
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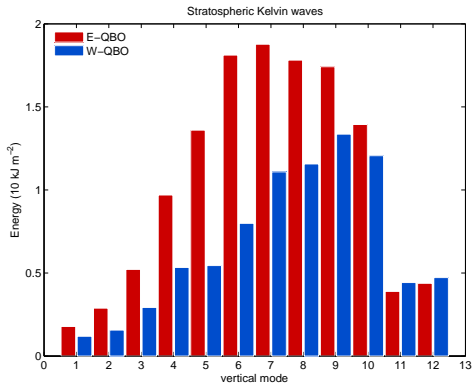
- ▶ Therefore the westerly QBO phase is associated with smaller vertical wavelengths.

Stratosphere : Ratio between the Normal Mode Power spectra for the two QBO phases (Easterly/Westerly)



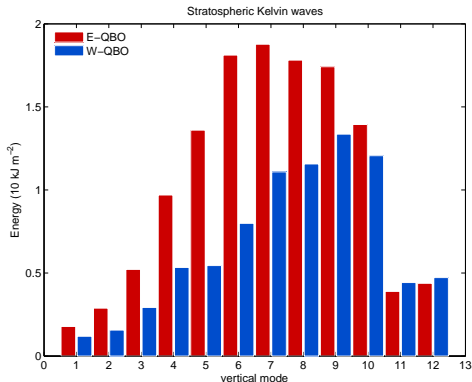
Stratospheric Kelvin waves





- Using typical values for the variables in the theoretical relationship

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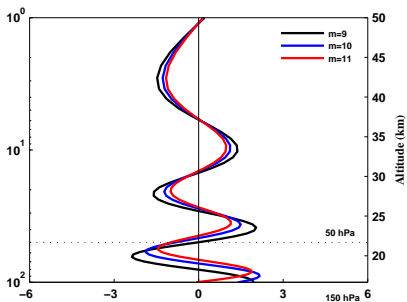
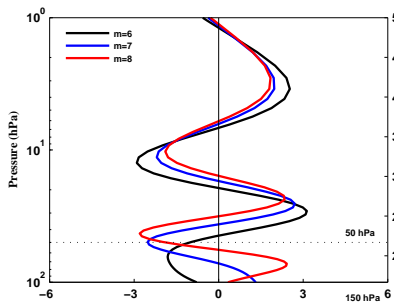
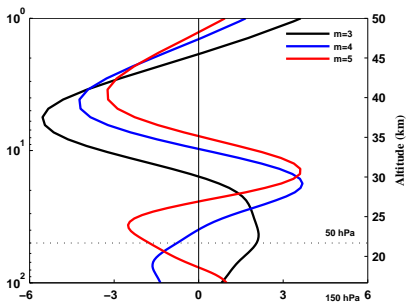
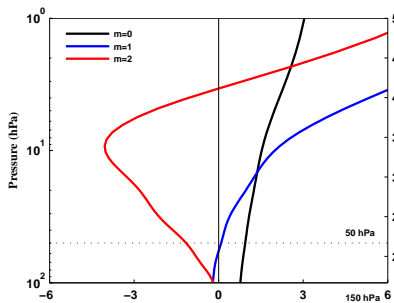
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- ▶ one obtains

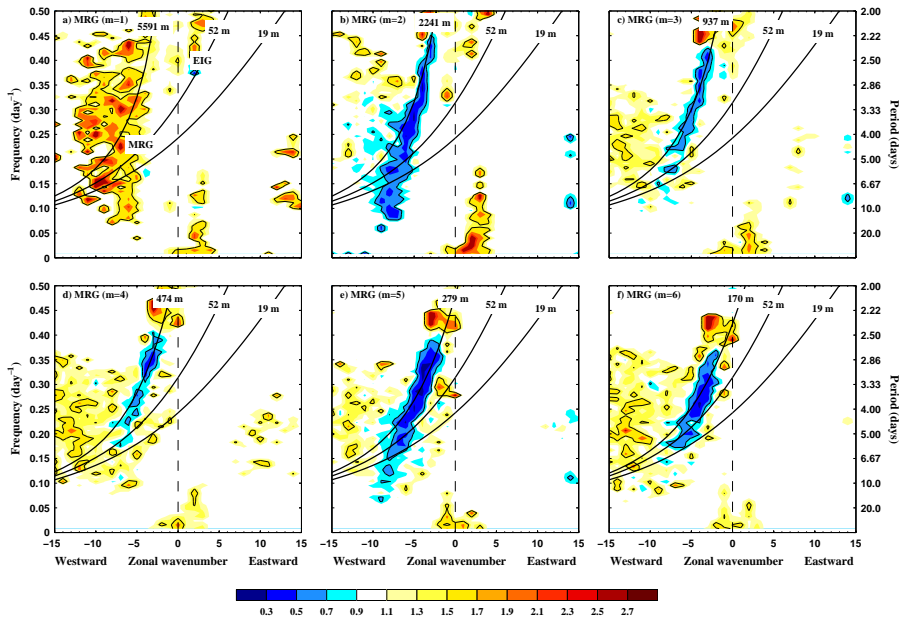
$$h_7 = 116 \text{ m} \quad \Rightarrow \quad L_z = 9.7 \text{ km}$$

$$h_9 = 52 \text{ m} \quad \Rightarrow \quad L_z = 6.5 \text{ km}$$

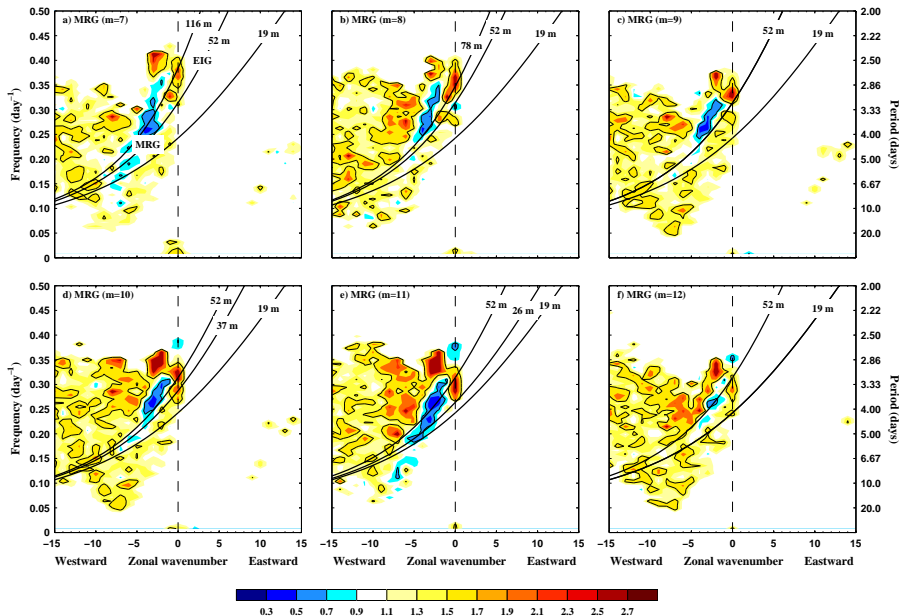
Vertical Structure Functions



Stratosphere : Ratio between the Normal Mode Power spectra for the two QBO phases (Easterly/Westerly)



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Thank you!