## Unstable periodic orbits and normal modes

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## **Barotropic atmospheric system**

The model is based on the barotropic vorticity equation on rotating sphere (=2D Navier-Stokes system +forcing +rotation +boundary&turbulent friction +orography).

$$\frac{\partial \Delta \psi}{\partial t} + J(\psi, \Delta \psi + l + H) = -\alpha \Delta \psi + \mu \Delta^2 \psi + f_{ext}.$$

 $\Delta$  - Laplacian, J - Jacobian, l - Coriolis parameter,  $\Psi$  - streamfunction, H - orography,  $f_{ext}$  - external (constant in time) forcing, turbulent viscosity  $\mu$ boundary layer friction  $\alpha$ , orography normalization k

Galerkin method, T21 resolution (spherical harmonics m<22), Phase space dim=231

## Model orography H



## External forcing f

$$f_{ext}=J(\psi_r,\Delta\psi_r+l+H)+\alpha\Delta\psi_r-\mu\Delta^2\psi_r.$$

 $\psi_r(t)$  is streamfunction on 300mb surface from 1960-1990 NCEP/NCAR reanalysis dataset



## **Parameters**

turbulent viscosity boundary layer friction  $\alpha = 2 \cdot 10^{-3}$  i.e. 1/25 days orography normalization k = 0.14

 $\mu = 6 \cdot 10^{-5}$ 

### **Model climates vs Nature**



300mb NCEP data

Average state

 $\overline{\psi} = \sum_{k} \psi(t_k) / K$ 

Variance

$$\sigma = \left(\frac{1}{K} \sum_{k} (\psi(t_k) - \overline{\psi})\right)^2)^{1/2}$$

T21

## **Model modes of variability**



Streamfunction EOF1,EOF2 EOF3,EOF4

## Attractor dimension equals to 62 28 positive Lyapunov exponents



Lyapunov spectrum for T21 models (Lyupunov time is 6 days)



## Linearization

# $\frac{\partial \Delta h}{\partial t} + J(h, \Delta \overline{\psi} + l + H) + J(\overline{\psi}, \Delta h) + \alpha \Delta h - \mu \Delta^2 h = 0$

 $\frac{\partial h}{\partial t} = Ah$ 





## Linear modes (streamfunction)

12 unstable modes (6 complex pairs), max(Re  $\Lambda$ ) + 1/14days Periods: 14, 8, 5, 56, 11, 31 days



#### EigVReal1, EigVCmplx1 EigVReal2, EigVCmplx2

Wavelike propagating structures with specific time-space scales Standard EOFs Bandpassed complex (rotational) EOFs

## **Complex (rotational) EOFs**

Consider system trajectory  $X_j(t^i)$  (mean removed). Conventional EOFs are e.v. of  $C = \langle XX^T \rangle$ 

- **1.** Apply direct (time) Fourier transform  $X_j \rightarrow F[X_j] = a_k + ib_k$  to every field component.
- 2. Apply  $\pi/2$  time shift
- **3. Apply inverse Fourier transform**
- 4. Make a complex time series and calculate its covariance matrix

$$a_k + ib_k \rightarrow b_k - ia_k$$

$$b_{k} - ia_{k} \rightarrow F^{-1}(b_{k} - ia_{k}) = \widetilde{X}$$
$$Z_{i} = X_{i} + i\widetilde{X}_{i}, C_{H} = \langle ZZ^{*} \rangle$$

5. Complex e.v. of  $C_H = \langle ZZ^* \rangle$  are Hilbert (complex, rotational) EOFs

## Leading Hilbert EOF of the system defines 2D oscillating structures with largest amplitude

#### E.V.1 (cmplx, real) (top)

VS

#### Complex EOF1 (9-16day bandpassed data) (bottom)



Pattern correlation 0.73/0.77

- E.V.1  $\rightarrow$  Complex EOF1 (09-16d bandpassed data)
- E.V.2  $\rightarrow$  Complex EOF1 (05-09d bandpassed data)
- E.V.3  $\rightarrow$  Complex EOF1 (02-06d bandpassed data)
- E.V.4  $\rightarrow$  Complex EOF2 (18-23d bandpassed data)
- E.V.5  $\rightarrow$  Complex EOF2 (09-16d bandpassed data)
- $\swarrow$  E.V.6  $\rightarrow$  Complex EOF1 (18-40d bandpassed data)



corr: (0.73, 0.77)

corr : (0.69, 0.67)

corr : (0.78, 0.78)

corr : (0.55, 0.44)

corr : (0.39, 0.55)

corr : (0.89, 0.60)

Normal modes of the system are strongly connected with variability patterns ! Why?

## **Linear approximation**

 $\frac{\partial \Delta \psi}{\partial t} + J(\psi, \Delta \psi + l + H) = -\alpha \Delta \psi + \mu \Delta^2 \psi + f_{ext}.$  $\frac{\partial h}{\partial t} = -\Delta^{-1}J(h,\Delta\overline{\psi} + l + H) - \Delta^{-1}J(\overline{\psi},\Delta h) - \alpha h + \mu\Delta h + \Delta^{-1}J(h,\Delta h)$  $\frac{\partial h}{\partial t} = Ah + R(h) \qquad R(h) \Longrightarrow -\lambda h + \zeta \\ A - \lambda E - \text{stable!}$  $\zeta$  - Gaussian w.noise  $\frac{\partial h}{\partial t} = (A - \lambda E)h + \zeta$ 

$$\frac{\partial h}{\partial t} = Lh + \zeta \quad L = A - \lambda E$$

PDF of the linear system is  $\rho(\psi) = c_o \exp((-C(0)\psi,\psi))$ 

#### 1. PDF must be gaussian



PDF of the T21 model projected at CEOF1 plane

## 2. Maximum density should be at the average state



Probability of the system trajectory to be in a ball around given state: average state (yellow), most visited trajectory point (green), average (black)

$$\frac{\partial h}{\partial t} = Lh + \zeta \quad L = A - \lambda E$$
$$h_{\Delta}(t) = \int_{t}^{t+\Delta} h(t') dt' / \Delta$$

Low frequency covariance matrix  $C_{\Lambda} = \langle h_{\Lambda}(t)h_{\Lambda}(t)' \rangle$ 

Uniform white noise assumption  $\langle \zeta(t)\zeta(t)' \rangle = k\delta(t-t')$ 

For linear-uniform white noise one can get  $\stackrel{\searrow}{\succ}$ 

 $C_{\Delta} \approx L^{-1} L^{-T} / k \quad \Delta \to \infty$ 

Low-frequency EOF are left singular vectors of inversed system linear operator



#### Leading left inv S.V. (top) / Leading low-freq. EOF (50+ d) (bottom)



First low-frequency EOF is reproduced but overall accuracy of "uniform white noise" approximation is not great.

#### Mean dynamics on leading CEOF1 plane (09-16d bandpassed)



Average tendency suggests rotation around mean state with period 12-20 days, rotation dominates far from the mean state. Nonlinearity (noise?) is not small.

 $T_{\delta} = 2\pi\delta / a.\sin\{\Psi_{12}'(\delta) \times \Psi_{12}') / |\Psi_{12}'(\delta) \| \Psi_{12}'|\}$ 

#### Mean dynamics on leading CEOF1 plane (05-09d bandpassed)



Average tendency suggests rotation around mean state with period 9-10 days, rotation dominates far from the mean state. Nonlinearity (noise) is not small.  $T_{\delta} = 2\pi \delta / a. \sin\{\Psi_{12}'(\delta) \times \Psi_{12}') / |\Psi_{12}'(\delta)| |\Psi_{12}'|\}$ 

## **Periodic orbits**

Periodic orbits are special trajectories of the system returning in the starting position after some T (period)

$$\psi(T) = S(T, \psi(0)) = \psi(0)$$

**1** Periodic orbits are the part of the system attractor and define recurrent (dynamical) circulation regimes.

2 Many chaotic systems have infinite number of periodic orbits.

**③** For axiom A and Anosov systems POs are dense on the attractor. Any system characteristic can be approximated by the set of the POs.

**4**Least unstable orbits are more important!

## **Methodology for finding orbits**

By the definition PO with period T is the system trajectory satisfying

$$\psi(T) = S(T, \psi(0)) = \psi(0)$$

This is the system of nonlinear equations having N+1 unknowns (initial point and period). Phase condition completes the system.

Damped Newton or inexact Quasi-Newton methods could be used to solve the system (+ line search, multi shooting, tenzor correction etc....)

•Initial guesses for u(0) and T - from the model trajectory { $\psi(k)$ , k=1..K} as local minimizers for  $|\psi(k+T) - \psi(k)|^2$ 

•What is required? Tangent model in full space (Newton) Or Tangent and adjoint tangent models in Krylov space (inexact methods). One iteration for finding UPO with period T means one run of tangent (adjoint tangent model) in full (Krylov) space for time T.

## **UPOs of the T21 system**

- **☆** 1500+ periodic orbits for T21 system were found
- All are unstable (25-60+ unstable directions), periods are 2-30+ days
- $\swarrow$  UPOs approximates system statistics as predicted!





#### Distribution of 800 most unstable UPOs by period



## UPOs on leading CEOF1 plane (09-16d bandpassed)



Spatial structure of leading UPOs matches structure characteristic time scale of CEOF1 and leading pair of normal modes.

#### **UPOs on leading CEOF1 plane (09-16d bandpassed)**



System shows regular rotational behavior when moving along CEOF1 plane. This is not always the case.

#### UPOs on leading CEOF1 plane (05-9d bandpassed)



Spatial structure of leading UPOs matches structure characteristic time scale of CEOF1 and second leading pair of normal modes.

## Summary

 Normal modes explain some patterns (leading CEOFs) of the system variability and corresponding timescales.



- It seems that linear approximation may not explain why.
  "Linear" behavior of the system occurs far from the mean state.
- 3. Least unstable POs of the system are important. Groups of UPOs are positioned along corresponding leading CEOFs (and normal modes). Periods of UPOs match normal mode frequencies. System moves along UPOs showing regular behavior.
- 4. Phase portrait of the system remembers its stable state analog?



$$h_{\Delta}(t) = \int_{t}^{t+\Delta} h(t')dt'/\Delta \quad C_{\Delta} = \langle h_{\Delta}(t)h_{\Delta}(t)' \rangle$$
$$C(\Delta) = \langle h(t+\Delta)h(t)' \rangle \quad \overline{C} = \int_{0}^{+\infty} C(\Delta)d\Delta$$
$$C_{\Delta} \approx (\overline{C} + \overline{C}^{T})/\Delta \quad \Delta \to \infty$$

Lyapunov equation

 $LC(0) + C(0)L^{T} = kE$   $C(\Delta) = \exp(L\Delta)C(0)$ 

 $C(0)L^{-T} + L^{-1}C(0) = kL^{-1}L^{-T}$  $\overline{C} = L^{-1}C(0)$  $\overline{C} + \overline{C} = kL^{-1}L^{-T}$